

Types of powder diffraction and high pressure experiments

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solid state matter = periodical arrangement of atoms
(**crystal**lographic structure → notion of **crystal**)

~100 different kinds of atoms
in the nature,
but
infinite ways to arrange these atoms



physical and chemical properties
of solid



precise knowledge of the crystallographic structure
of a material to understand these properties.

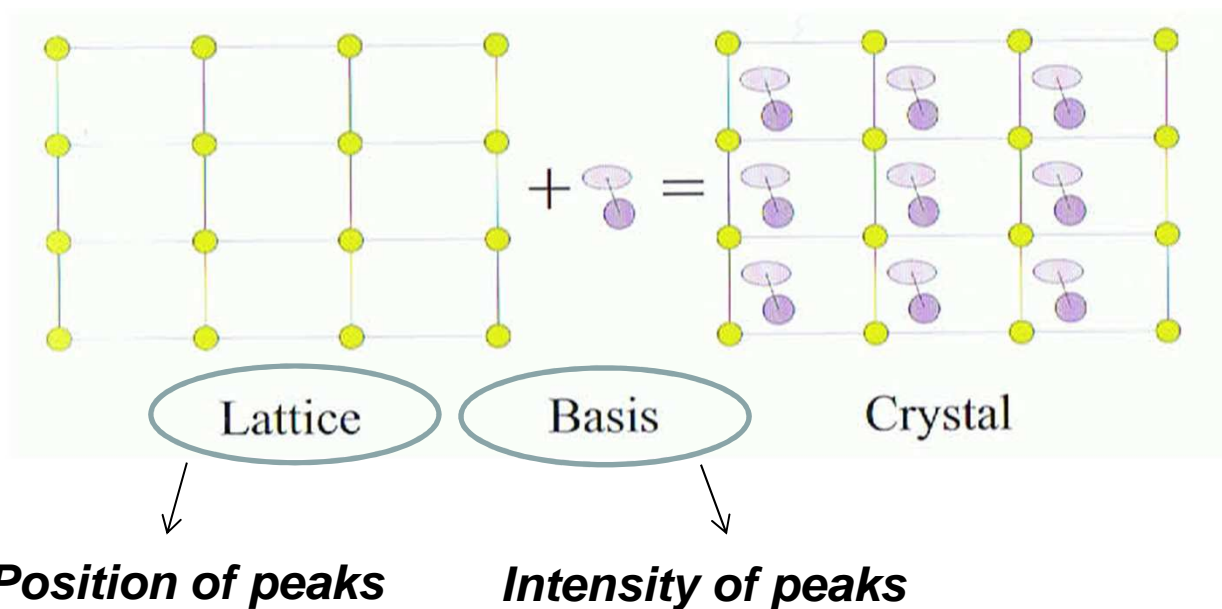


The elastic scattering of a radiation beam (diffraction) represents the
most suited technique to determine crystallographic structures
(nowadays widely used in chemistry, physics, engineering, and biology).

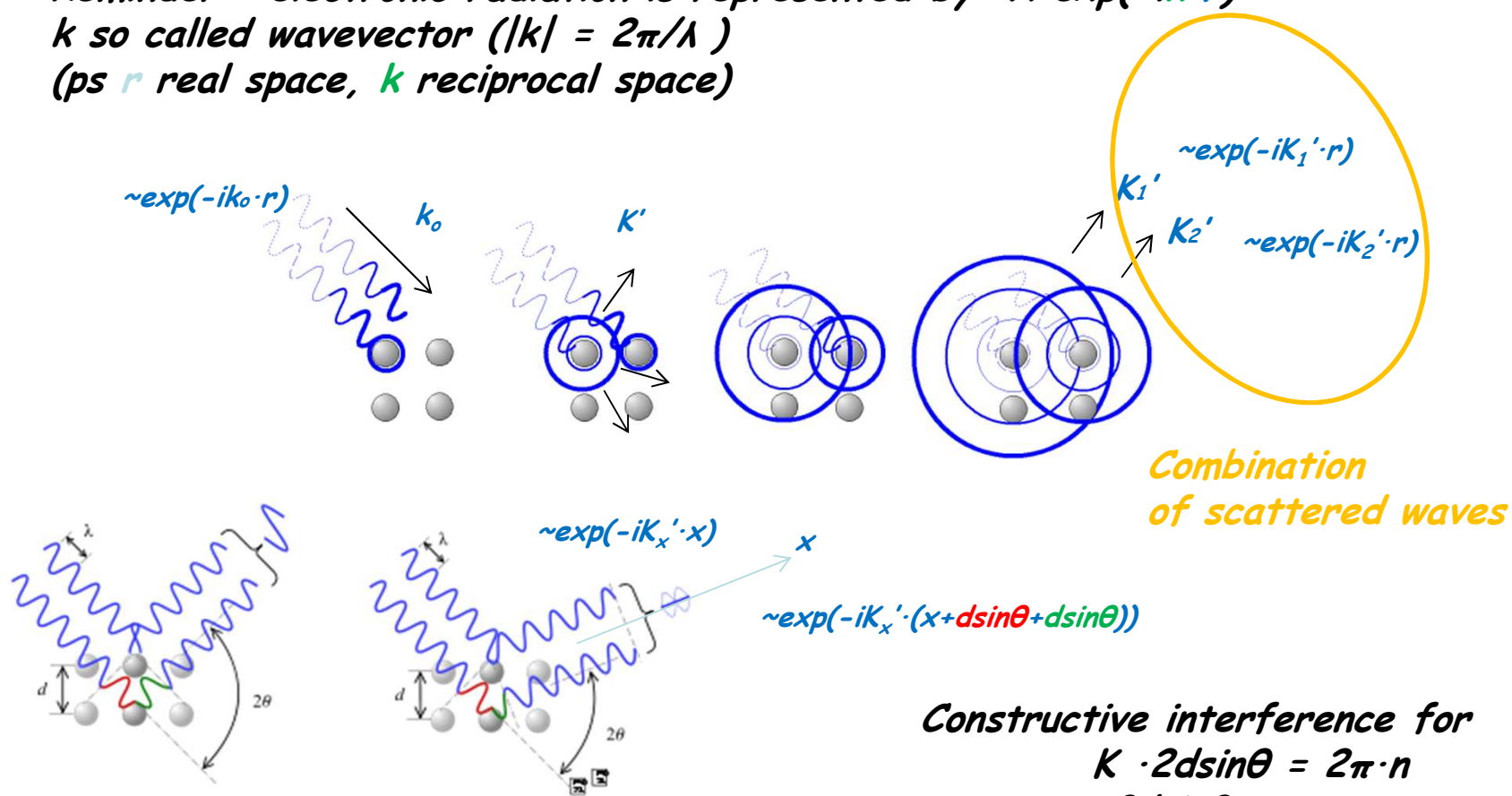
What is a crystal ?

- Well it is neither a material in liquid, nor in gaseous state -> not amorph
- So it is a solid

A crystal is a 3-dim periodic (and ideally infinite) arrangement (**lattice**) of identical scatterers (**basis**). (basis = molecule)
Each point has an identical surrounding



Reminder : electronic radiation is represented by $A \exp(-ik \cdot r)$
 k so called wavevector ($|k| = 2\pi/\lambda$)
 (ps r real space, k reciprocal space)





Bragg formulation of X-ray diffraction

“crystals, at certain specific wavelengths and incident angles, produce intense peaks of reflected radiation (known as Bragg peaks)”

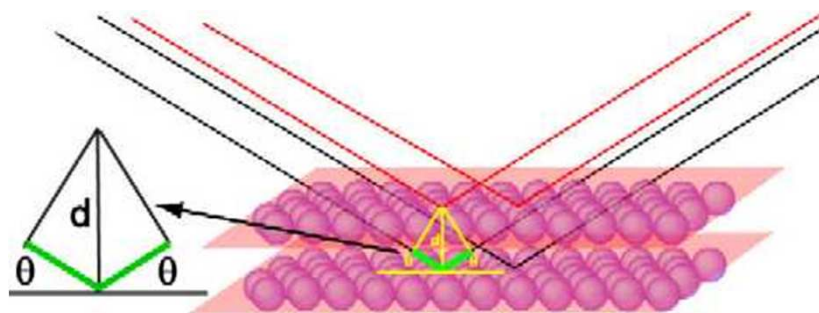
“W. L. Bragg explained this result by modeling the crystal as a set of discrete parallel planes separated by a constant parameter d ”

$$\lambda = 2d \sin \theta$$

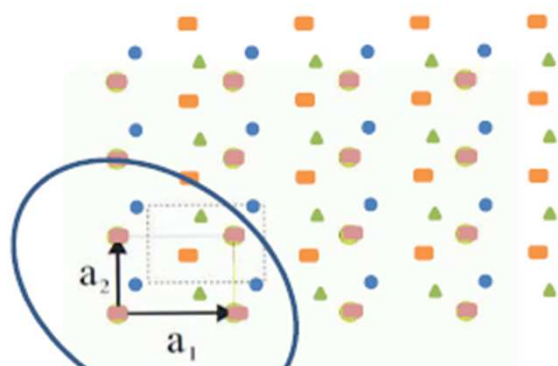
d spacing between *planes*
 λ *wavelength* of radiation

d, λ must have same
order of magnitudes

planes are formed by layer of atoms $\rightarrow \sim 1\text{-}100 \text{ \AA}$



*Typical *wavelength* typical wavelength of X-rays or neutrons $\rightarrow \sim 1\text{-}100 \text{ \AA}$*



In unit cell system,
the 4 atoms are at

- $x_0=0, y_0=0$
- x_1, y_1
- x_2, y_2
- ▲ x_3, y_3

$$F_{hkl} = f_0 \cdot \exp(iQ_{hkl} \cdot r_0) + f_1 \cdot \exp(iQ_{hkl} \cdot r_1) + f_2 \cdot \exp(iQ_{hkl} \cdot r_2) + f_3 \cdot \exp(iQ_{hkl} \cdot r_3)$$

The intensity of a specific Bragg reflection hkl is proportional to the so called structure factor

$$F(Q_{hkl}) = \sum_j f_j e^{iQ_{hkl} \cdot r_j} t_j(Q_{hkl})$$

summation over
all atoms in the unit
cell

temperature factor

position of each atom
in unit cell

scattering strength power of the atom (radiation and
energy dependent)

$$\text{X-ray: } f_j = f_0 + f'(E) + i f''(E)$$

f_0 = number of electrons

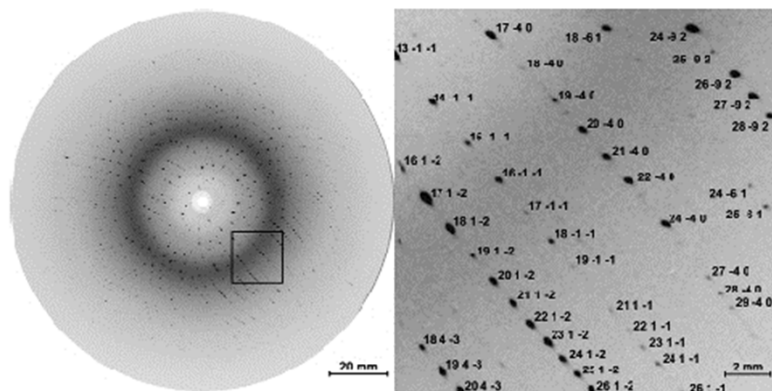
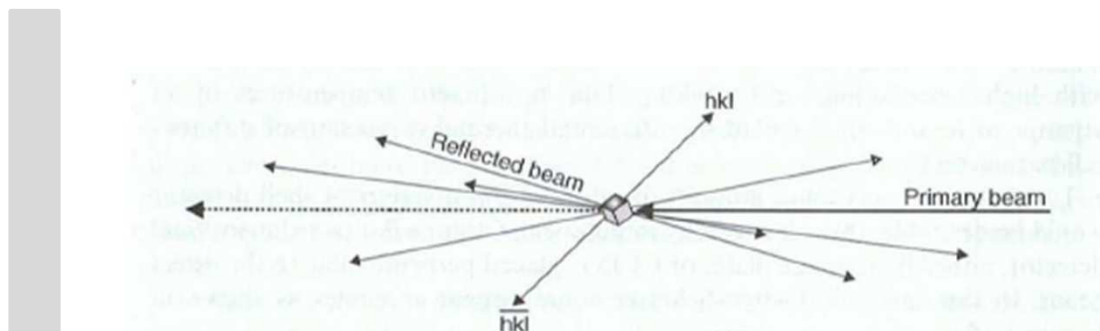
$f', f'' \neq 0$ around absorption edge

$$\text{Neutrons: } f_j = b$$

b no systematics vs. periodic table, it is
related to atom nucleus

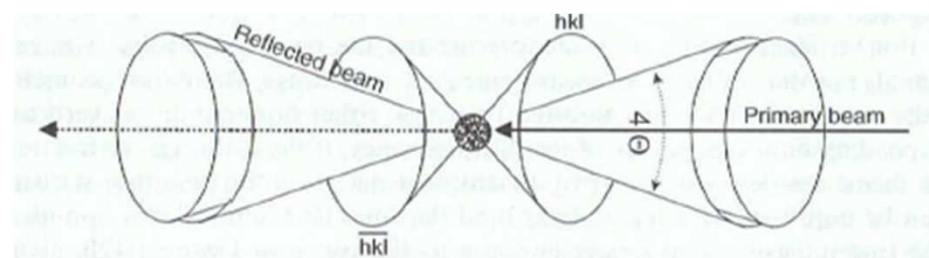
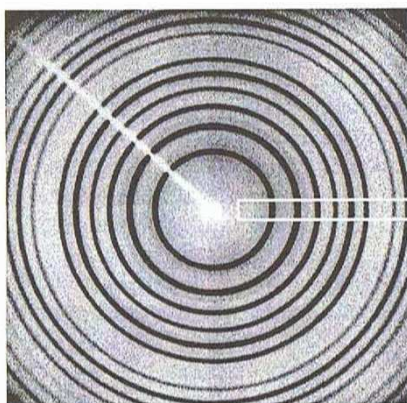
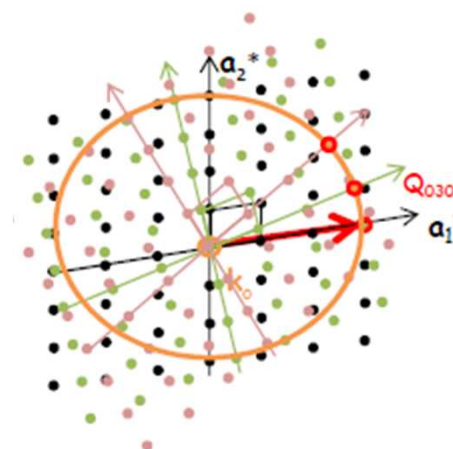
b can be related to the intrinsic magnetic
moment of atoms

detector



*Rotate the crystal to get other
Bragg reflections
Complete experiment :
100-300 frames at various crystal
orientation*

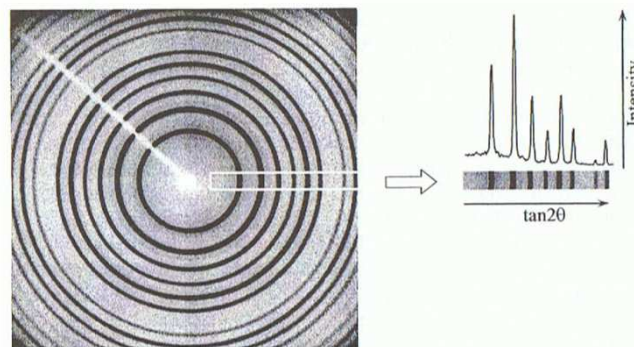
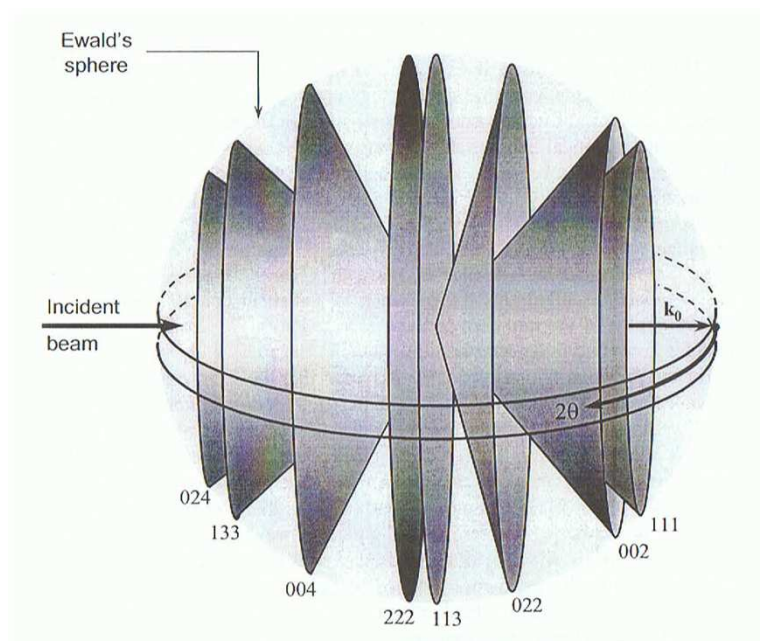
- A powder is a multitude of single crystals, the diffraction by a powder is then the superposition of all contributing crystals.
(a powder grain is not necessarily a single crystal)
- The hkl bragg peak on the detector is no longer a single spot in space but the sum of spots arising from each contributing crystal *(problematic of statistical averaging)*
- All the hkl spots are at 2θ angle from incoming beam, they hence all lie on the so called Debye-Scherrer cone.
- Intercept of detector plane with cone produces rings

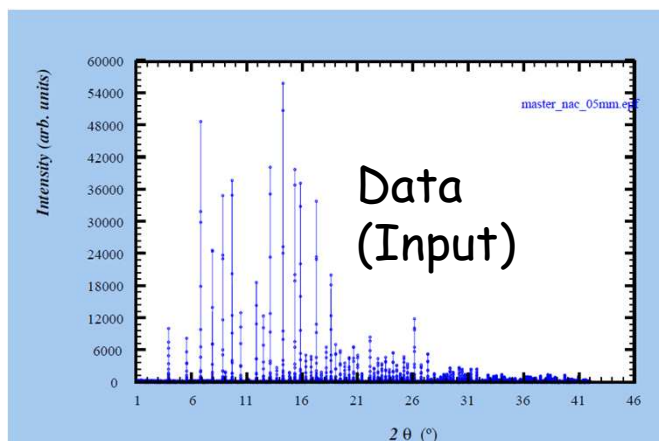


Debye-Scherrer cone

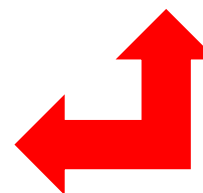
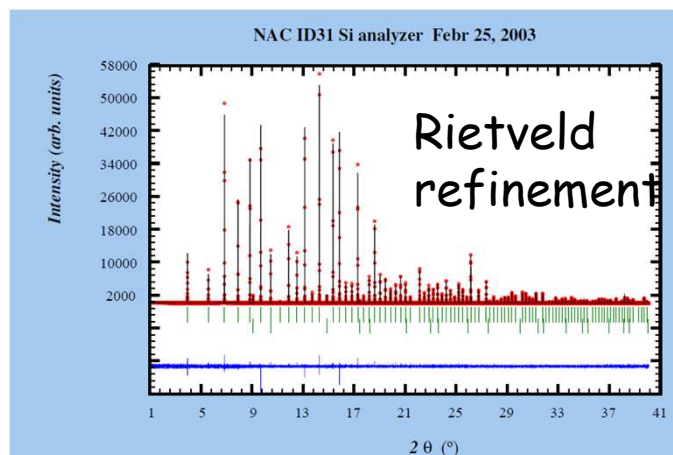
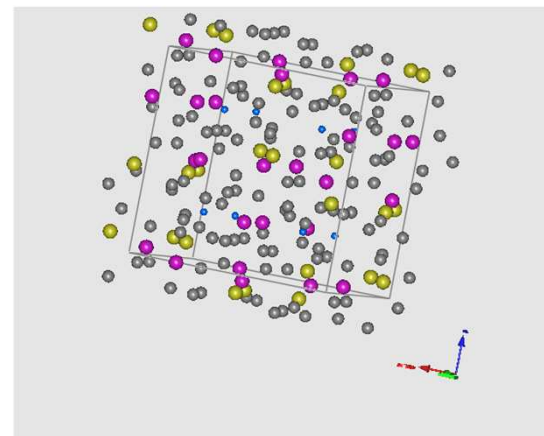
A diffraction experiment consists in

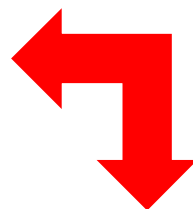
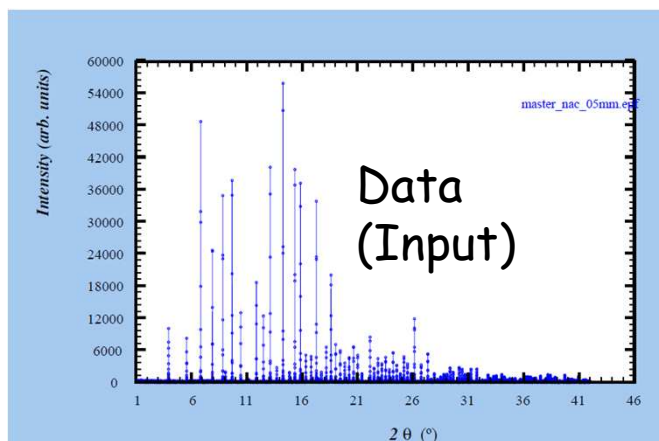
- 1) setting a crystal or a powder (= statistical ensemble of crystals with "random" orientation) in a well known radiation beam*
- 2) measure the (non-isotrop) level of radiation scattered by this material.*





Model (Input/output)

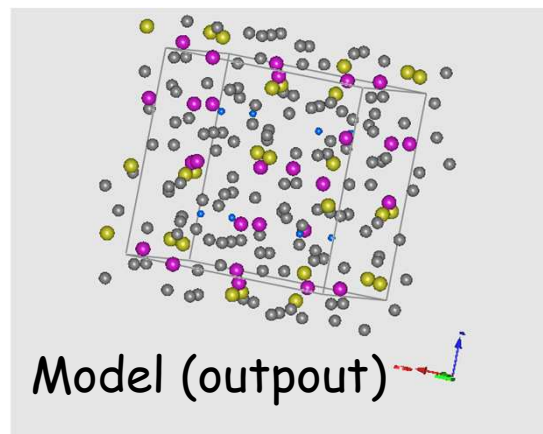


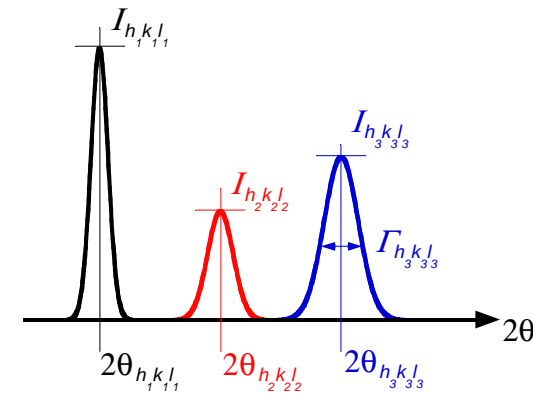
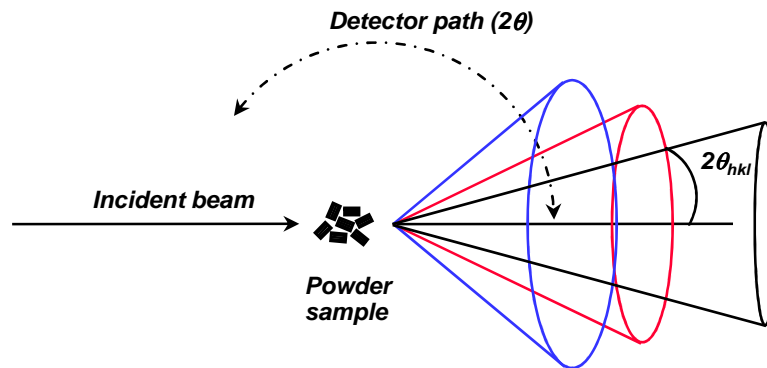


- Assign hkl for each of the peak
→ (extract space group)
- extract integrated intensity
of each individual peak
→ extract structure factors

Alternative to Rietveld ?

So called ab-initio or direct methods





The center θ_g

Bragg law :

$$\lambda = 2 \cdot d_{hkl} \cdot \sin \theta_{hkl}$$

$\Rightarrow d_{hkl} = d_{hkl}(a, b, c)$

\Rightarrow information on
overall periodic
arrangement of atoms

The intensity $I_g = f(\theta_g)$

Structure factor:

$$F_{hkl} = \sum_j f_j \exp[-i2\pi(x_j h + y_j k + z_j l)]$$

\Rightarrow kind of atom (f_j) in material

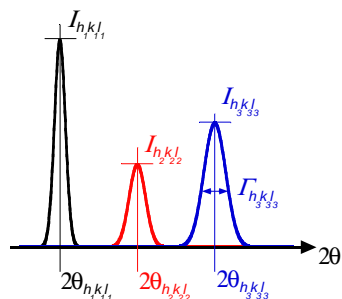
\Rightarrow position (x_j, y_j, z_j) of atoms
in the structure (unit cell)

The fwhm Γ_g

$$(f(\theta_g + \Gamma_g) = 1/2 \cdot f(\theta_g))$$

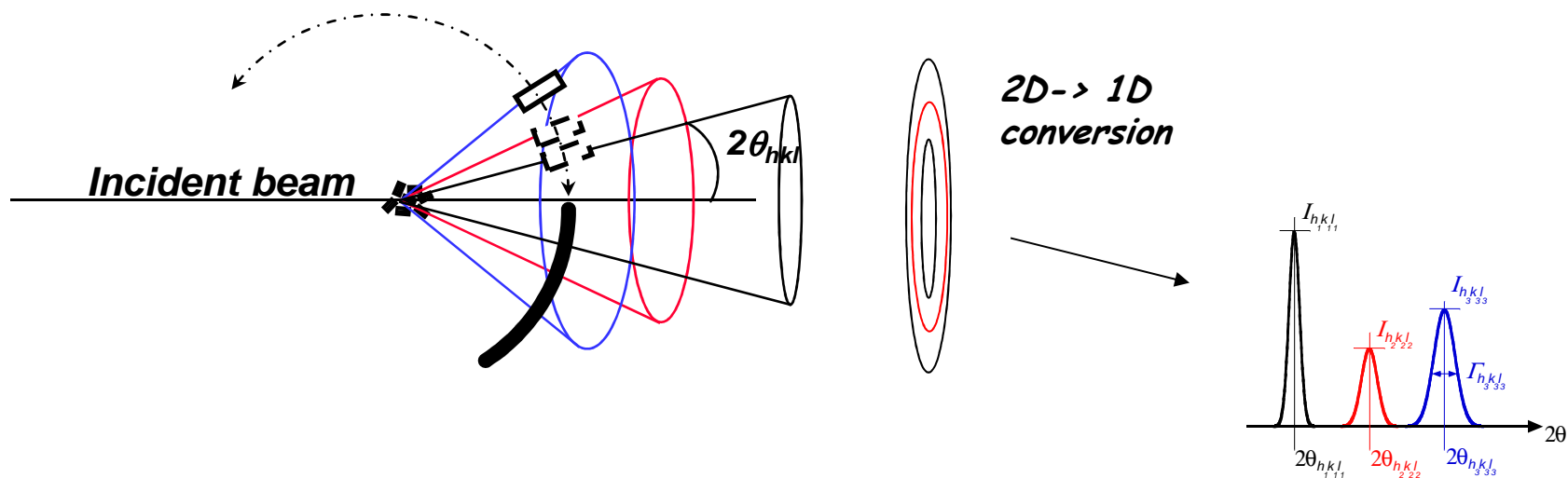
\Rightarrow microstructural
features (strain, size)

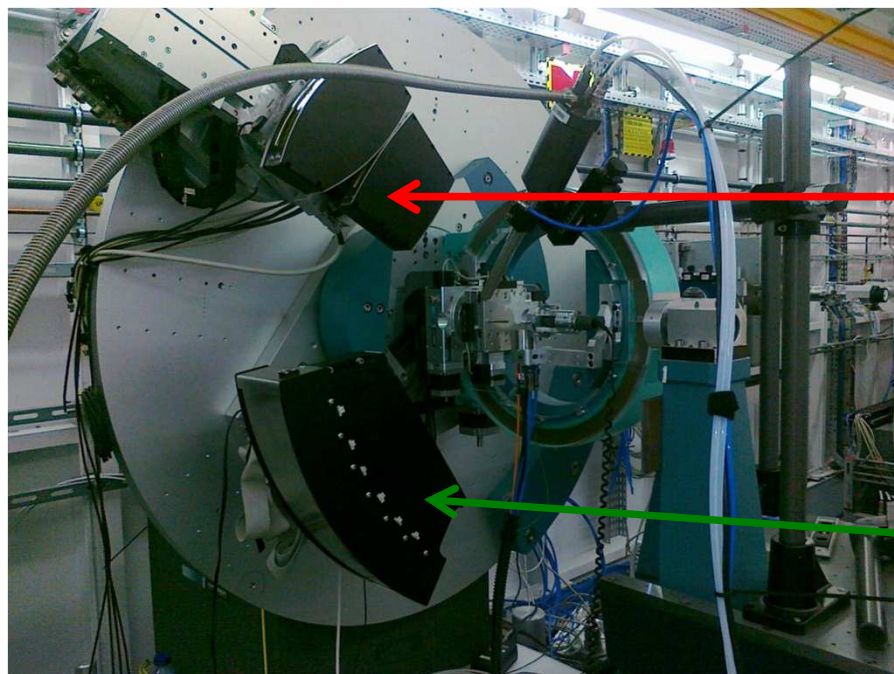
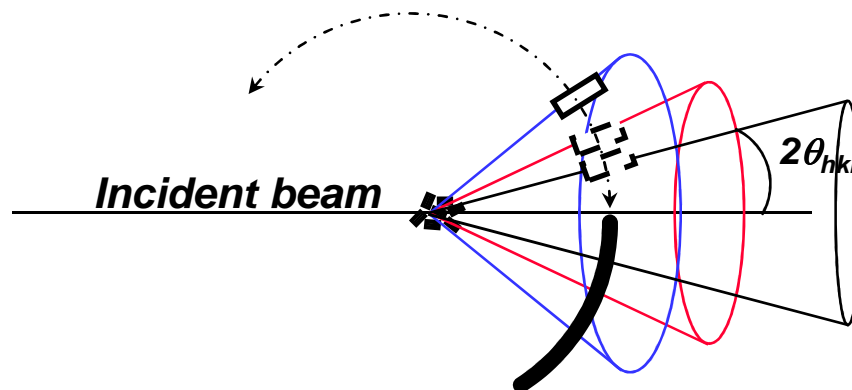
strongly instrument
dependent



direct 1D acquisition
 either scanning a point detector (0D)
 or collecting on a 1D steady state detector

Key elements :
 the source, the detectors
 the mechanics holding the sample and detector



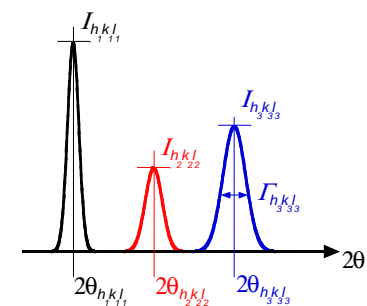
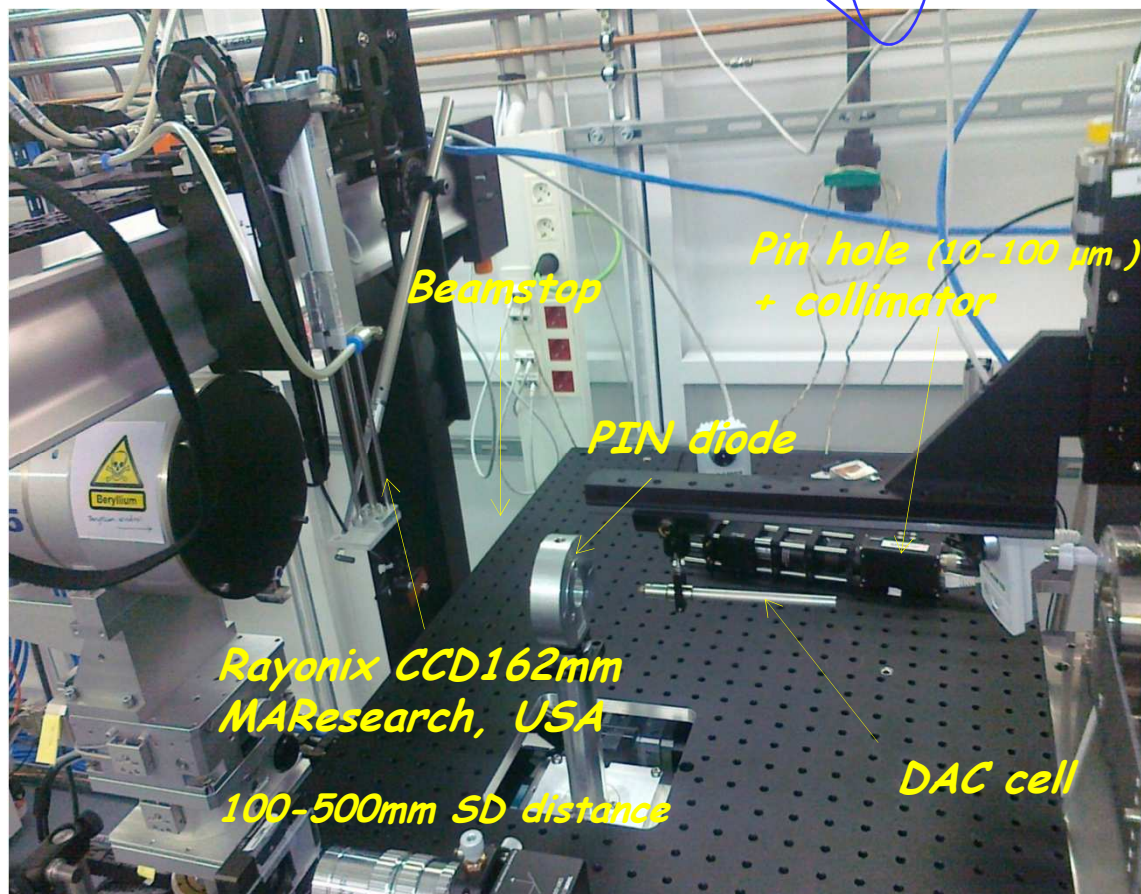
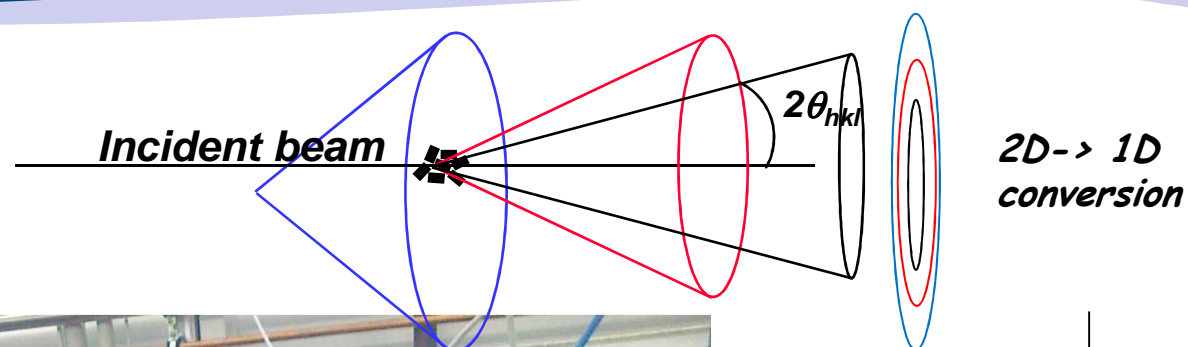


Multi Analyzer Detector (10? - 50 keV)
(to be commissioned Nov 2012)

- 13 channels with 1.5 deg pitch
- Si_{111} or Si_{220} Bragg reflection
- YAP scintillator + PMT

Mythen detector (7-30? keV)

- 6 modules (1280 channels, 50 μ m pitch)
- D_{SD} = 550mm
- ms time resolution
- -> ~40deg in 0.005 deg pitch

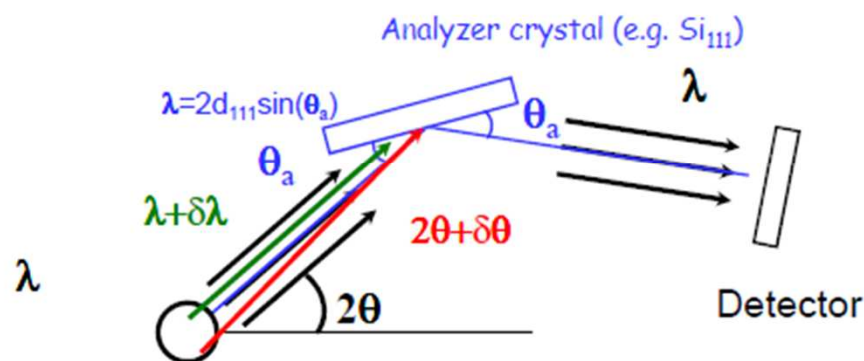


How to achieve higher resolution better signal to noise

1- minimize incoming beam divergence and energy resolution

(no focusing, collimating mirror, DCM monochromator, synchrotron source
(eg. Wiggler => high horizontal divergence, low vertical

2- use analyzers (back monochromator)

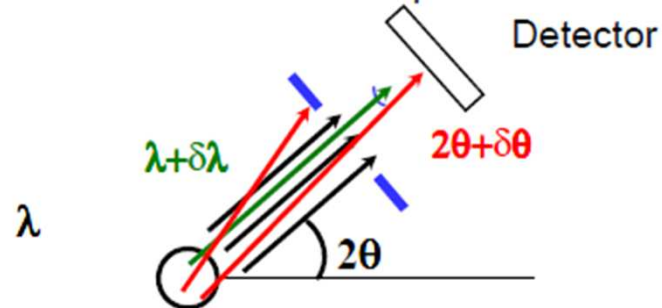


high signal to noise ratio
(elimination of inelastic processes)

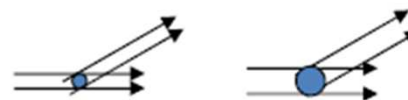
higher angular resolution
(analyzer crystal acts as a tiny slit:
 $d\theta \sim 10'' \sim$ intrinsic width of the
analyzer crystal)

no sample size/transparency effect
Not sensitive to displacement

Slit/No Slit-detector setup



high intensity rate
(since no loss (discrimination)
from analyzer-radiation interaction)
Angular resolution is proportional to
sample size in the beam

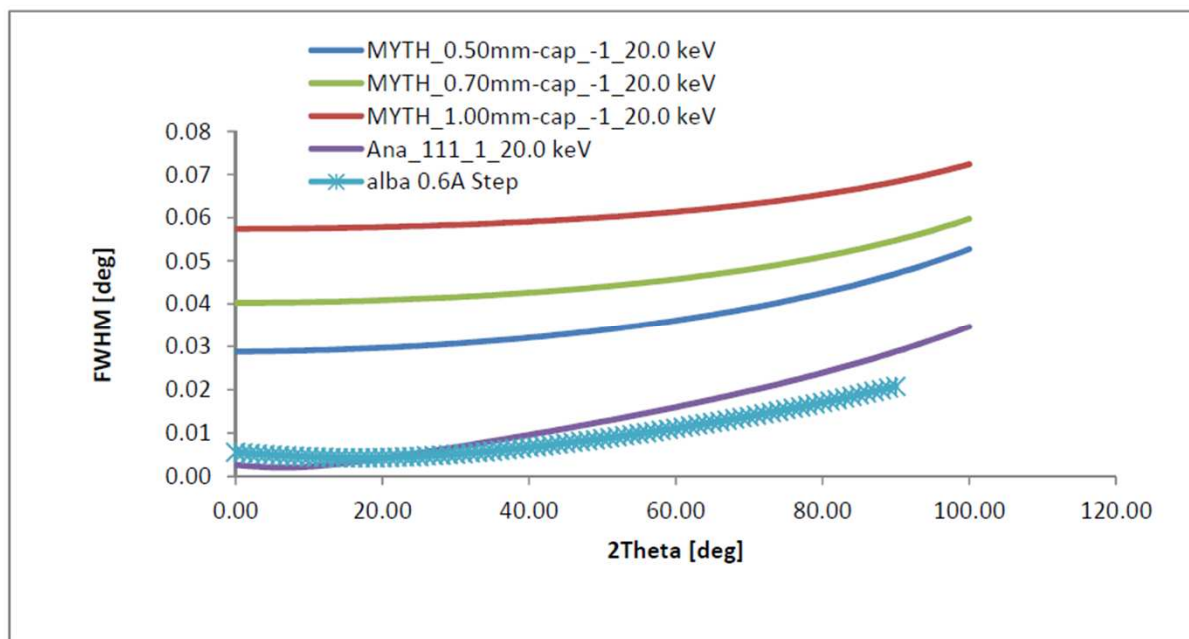


$$\Gamma_{ana} = \sqrt{\omega_a + \frac{1}{2} \omega_m \left\{ \frac{\tan(\theta_a)}{\tan(\theta_m)} - 2 \frac{\tan(\theta_s)}{\tan(\theta_m)} \right\}^2 + \varphi_{M1}^2 \left\{ \frac{\tan(\theta_a)}{\tan(\theta_m)} - 2 \frac{\tan(\theta_s)}{\tan(\theta_m)} \pm 1 \right\}^2 + \varphi_{M2}^2}$$

$$\Gamma_{mythen} = \sqrt{\sqrt{d^2 + p^2} + \frac{1}{2} \omega_m \left\{ -2 \frac{\tan(\theta_s)}{\tan(\theta_m)} \right\}^2 + \varphi_{M1}^2 \left\{ -2 \frac{\tan(\theta_s)}{\tan(\theta_m)} \pm 1 \right\}^2 + \varphi_{M2}^2}$$

$\omega_{a/m}$ Darwin width mono/analyzer Xtal
 $\varphi_{M1/M2}$ Beam divergence on before/after
 $\theta_{a/m/s}$ Angle sana/mono/sample (energy)
 p, d Pitch between stripe, capillar diamteer

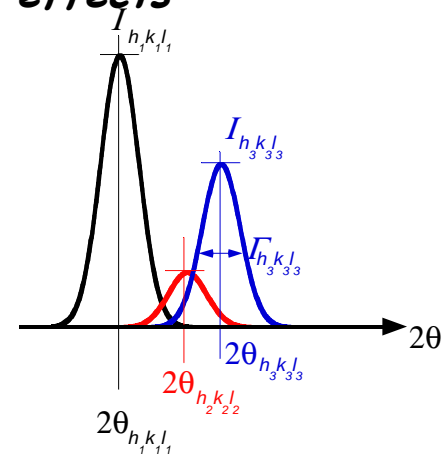
+ additional horizontal divergence
and sample size contribution



Why high resolution and high accuracy ?

*Not only to get accuracy of cell parameter up to forth or fifth digit
As seen before, in order to detect microstructural effects
But as well to minimize the overlap problem*

In some circonstances (short λ , high 2θ , big unit cells) Bragg reflections strongly overlap



*characteristic values
(θ_g , I_g , Γ_g) difficult to extract*

*inaccuracy in the
determination of a crystal
structure*

Reflections in a powder pattern can be described by various profile functions

The most important are

- **Gaussian** function $G(x)$
- **Lorentz** (or Cauchy) function $L(x)$
- **pseudo Voigt** function $pV(x)$

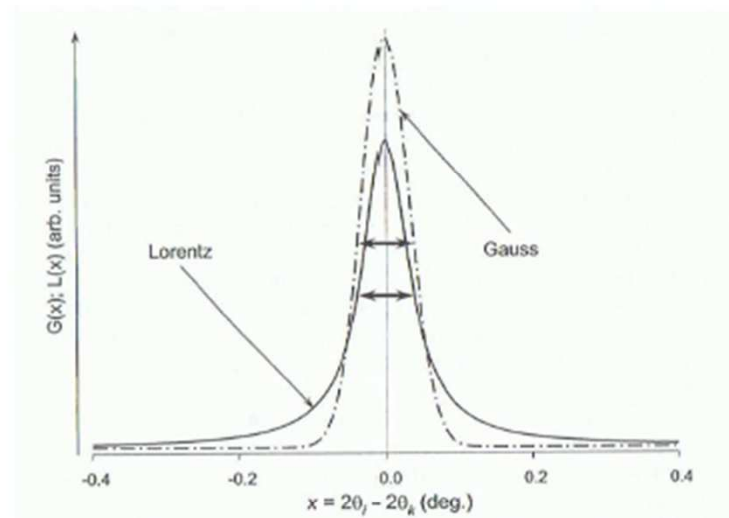
$$G(x) = \frac{\sqrt{C_G}}{\sqrt{\pi}H} e^{-C_G x^2} \quad L(x) = \frac{\sqrt{C_L}}{\pi H} (1 + C_L x^2)^{-1}$$

$$C_G = 4 \ln 2$$

$$C_L = 4$$

$$pV(x) = \eta L(x) + (1 - \eta)G(x)$$

H : full width at half maximum FWHM $pV(x \pm H/2) = \frac{1}{2} pV(x)$



Why is the easy modelisation of shape function important ?

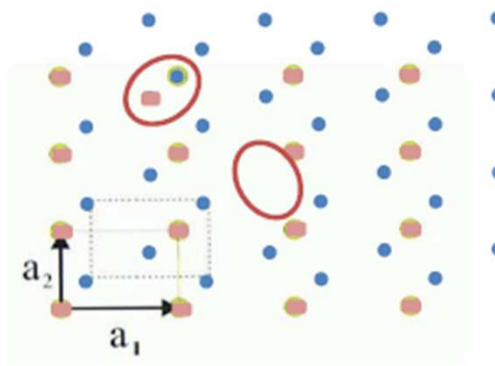
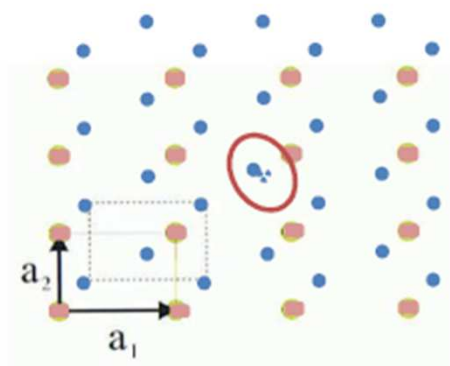
Data analysis of powder patterns by so called Rietveld refinement implies comparison of the measured spectrum with a trial calculated pattern. This process involves calculation of thousands of Bragg reflection.

Rapide calculation only possible as long as Bragg reflection are easily modeled.

Why is the knowledge of shape function important ?

*Width Γ_g information on microstructural features,
that is any form of atomic disorder or imperfections exhibited by the
periodical arrangement of atoms
(crystallite-size effects, dislocations, stacking fault, variations in interatomic
distances due to internal stresses or non-stoichiometry)*

-> all what contributes to reduce the infinite periodicity of the lattice

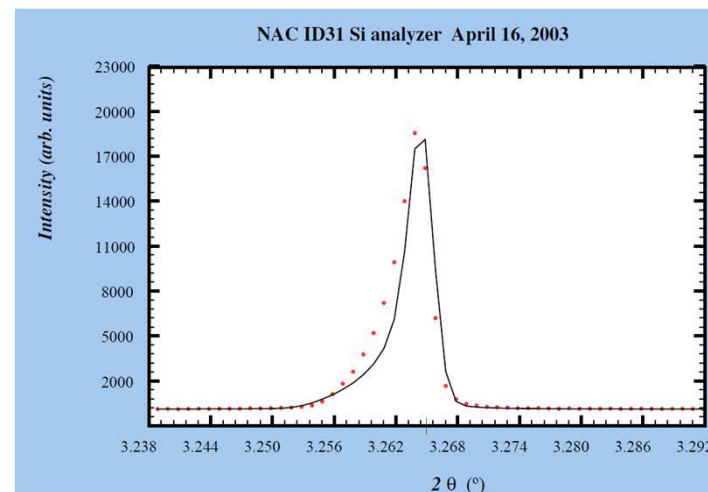
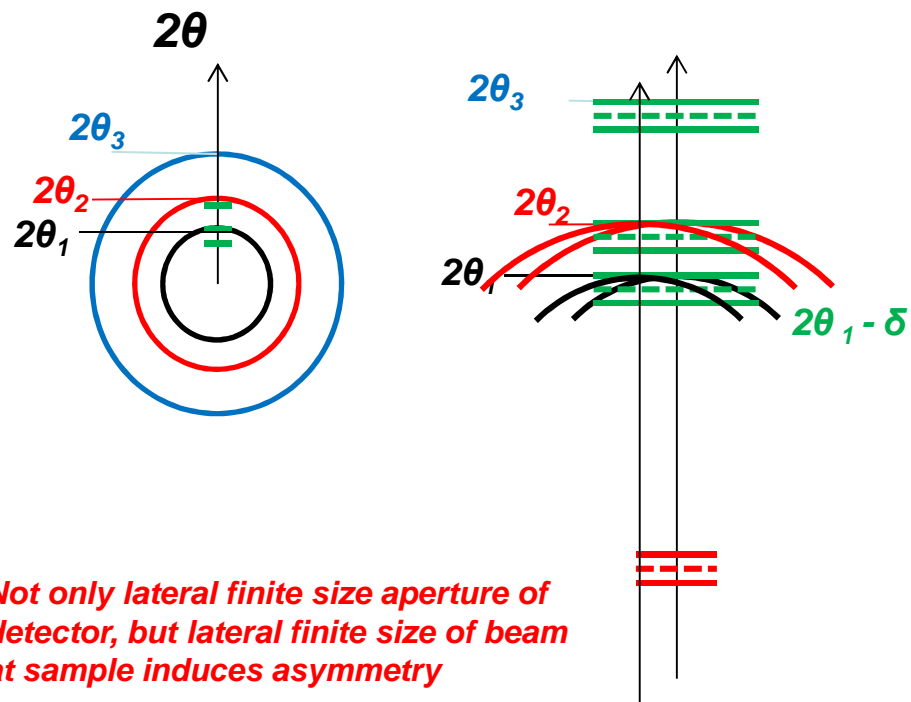


■ *mixed site*

*2-3 types of atoms
(generally of the
same size), occupy
the same lattice site*

*These features are very small and generally contribute marginally to the width.
The main contribution comes from the instrument itself -> high resolution needed*

*Sever peak shape asymmetry at low angles
 -> due to limited size aperture of detector and beam*



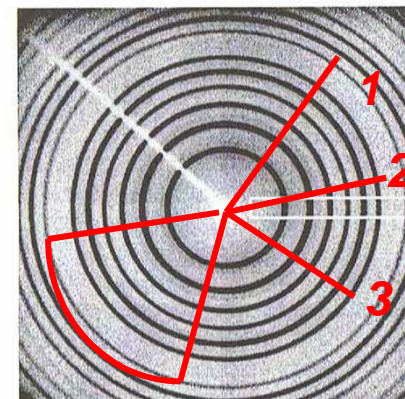
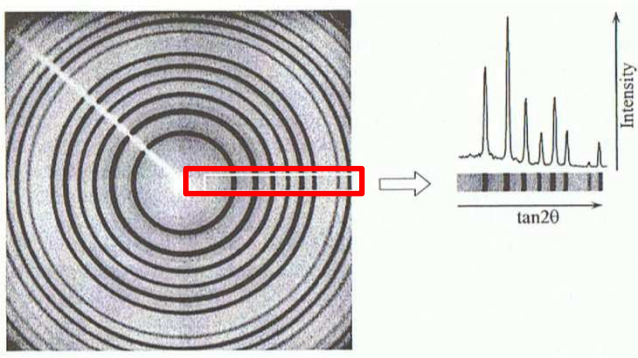
**So called “umbrella” effect can
 Be ‘hardly’ corrected using so called
 Thomson-Cox-Hasting asymmetry
 correction function**

**depends on S/D and L/D where
 S lateral dimension of beam at sample
 L lateral dimension of detector aperture
 D sample to detector distance**

How to get rid of asymmetry

- Minimize S and/or $L \rightarrow$ lost of intensity
- Increase Sample to deetctor distance $D \rightarrow$ lost of intensity
- Use higher wavelength so to have Bragg peaks at higher 2θ

No asymmetry if using 2D detectors and properly convert from 2D to 1D



1+2+3

Why using Synchrotron ? (compared to Lab X-ray sources)

- *Energy tunability (not limited Cu K, Mo K edges)*
- *Broader range of energies (8-50 keV at MSPD@ALBA)*
- *High energies = higher penetration depth:*
 - > *transmission geometrie*
 - > *sample environments*
(cryostats, furnace, ...all what includes windows)
- *High flux*
 - > *high throughput*
 - > *high resolution*
 - *we can afford putting "beam cleaning elements*
(DCM monochromator, analyzer crystals)
 - *very low natural beam divergence*
 - > *small spot size*
- *Big diffractometers*
 - *Space to put ancillary equipments*

Why using Synchrotron ?

Practical examples

- *When good signal to noise is needed*
- *When high angular resolution is needed -> resolve peaks*
- *When energy tunability is needed -> resonant diffraction*
- *When high throughput is useful*
- *When temperature helps*
- *Beyond the reasonable*
- *When small spot size is needed -> high pressure*
- *When high energy and controlled spot size is needed*
- *When high angular resolution is needed -> microstructural effects*

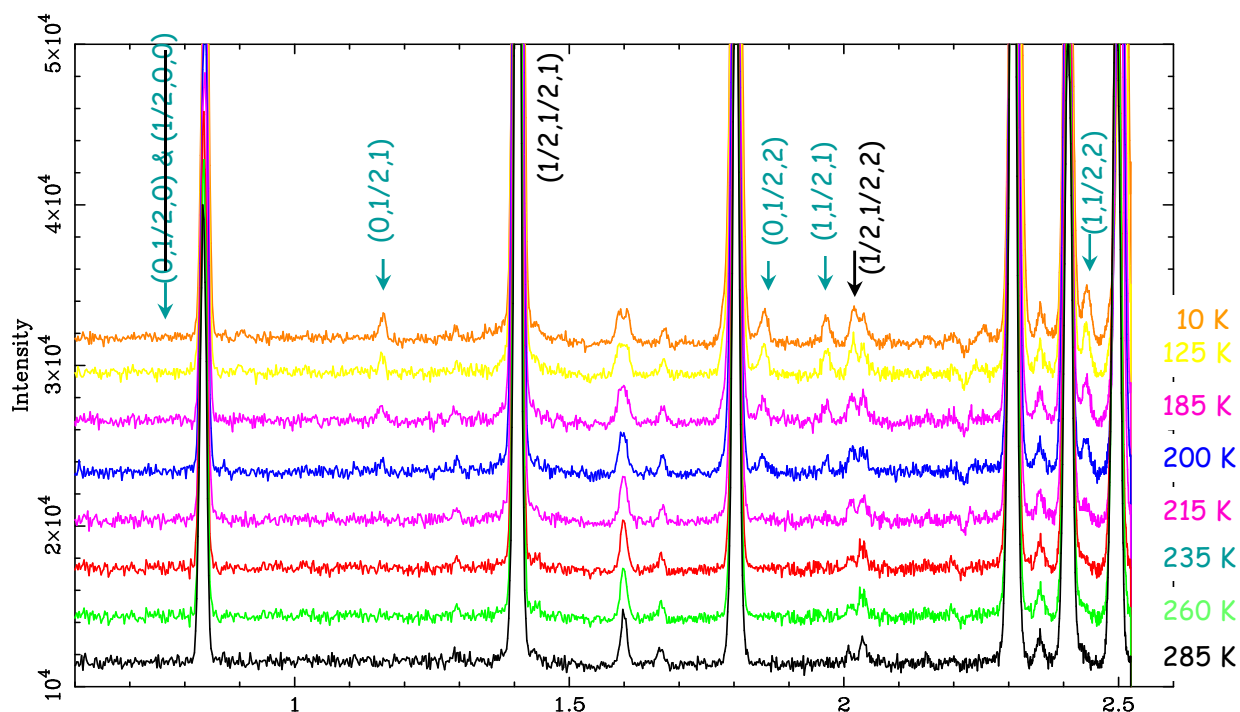
$\text{LnBaCo}_2\text{O}_{5.0}$ ($\text{Ln}=\text{Ho}^{1,2}, \text{Tb}^2, \text{Dy}^2, \text{Y}^3$) Charge Ordering

- 1) Suard et al., Phys. Rev. B61 (2000) R11871
 2) Fauth et al., Eur. Phys. J B21 (2001) 163
 3) Vogt et al., Phys. Rev. Lett. 84 (2000) 2969

Neutron Powder Diffraction : G61@LLB ($\lambda=4.73 \text{ \AA}$)

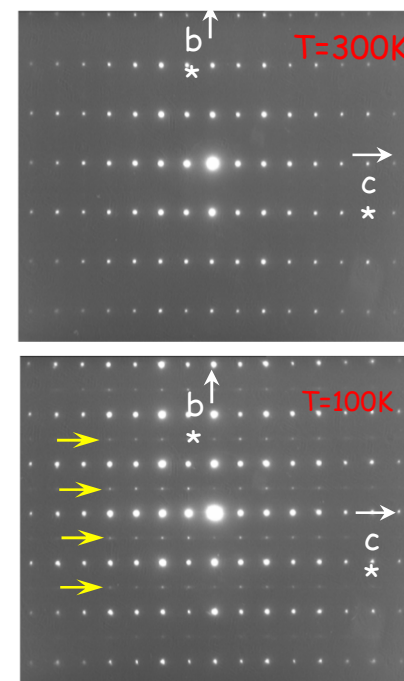
$\text{HoBaCo}_2\text{O}_5$: (indexation in $a_p \times a_p \times 2a_p$ setting)

Standard Electron diffraction
 $\text{TbBaCo}_2\text{O}_5$ (+ microdiffraction)



no $(0,1/2,0) \rightarrow \text{S.G. Pmmb}$ $Q [\text{\AA}^{-1}]$

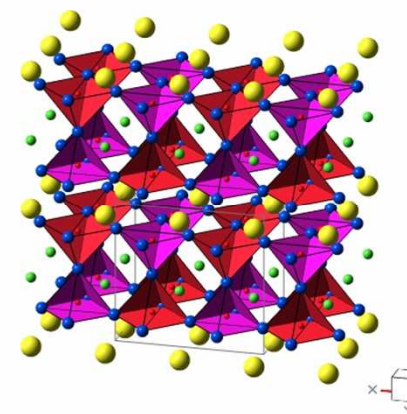
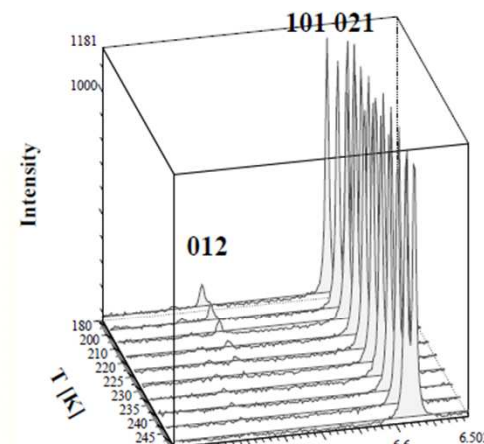
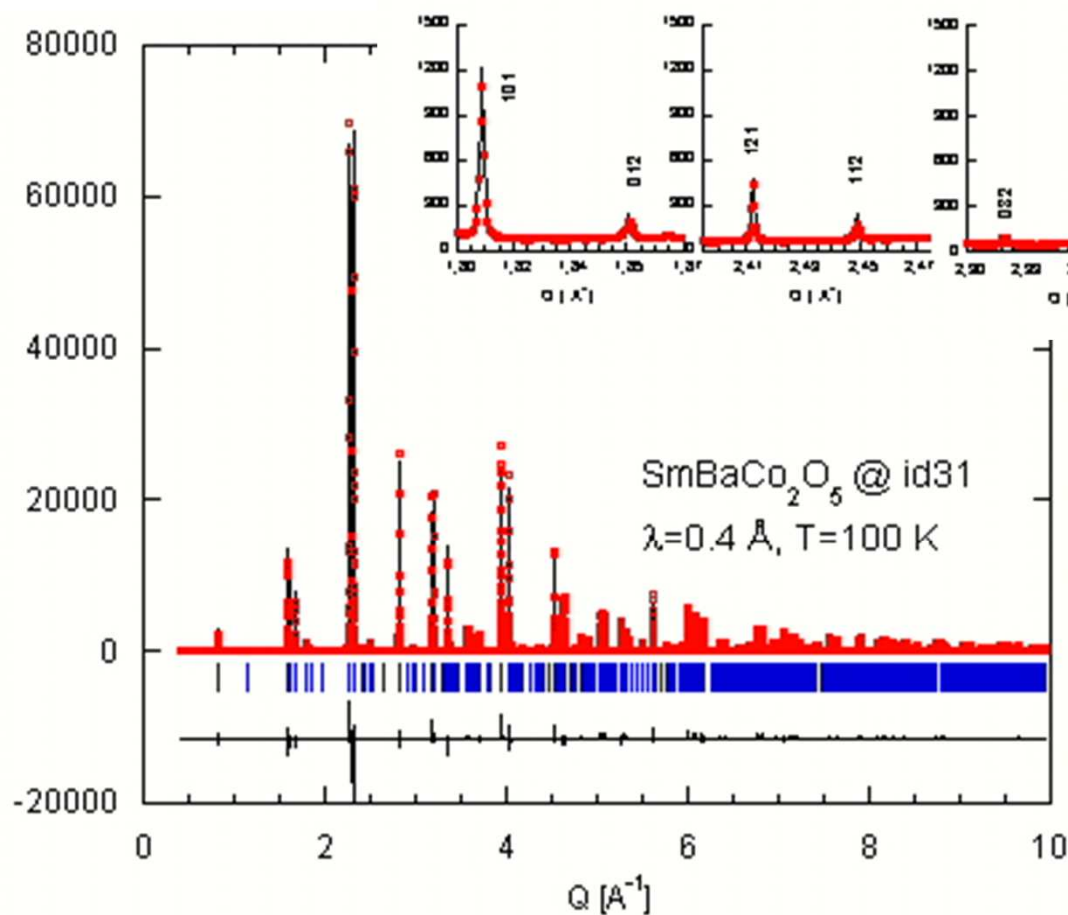
no $(1/2,0,0) \Rightarrow$ magnetic moments along the a -axis



Pmmb extinction
 $0k0, hk0, k=2n+1$

Charge Ordering in $\text{Ln}=\text{Y}, \text{Ho}, \text{Tb}, \text{Dy}$ and $\text{SmBaCo}_2\text{O}_{5.0}$

Synchrotron Powder Diffraction :
Pmmn Co structure



What about CO in $\text{LnBaCo}_2\text{O}_5$ $\text{Ln}=\text{Nd}, \text{Pr}$?

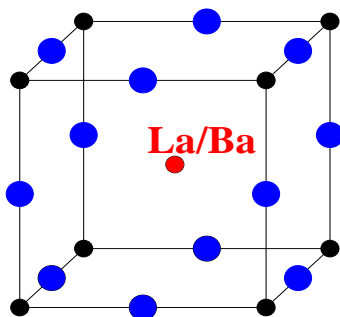
What about CO in $\text{LnBaCo}_2\text{O}_{5+e}$? (e_{max} ?)

LaBaCo₂O_{6.0}¹ : Orbital Ordering

1) Fauth et al., Phys. Rev. B65 (2002) 060401(R)

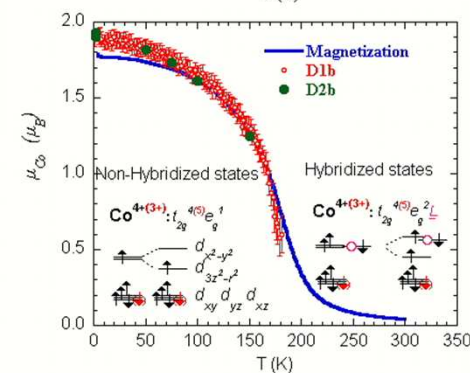
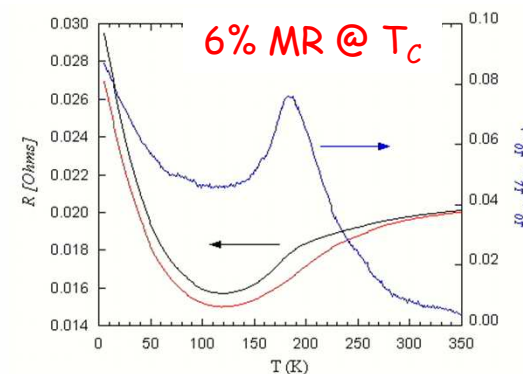
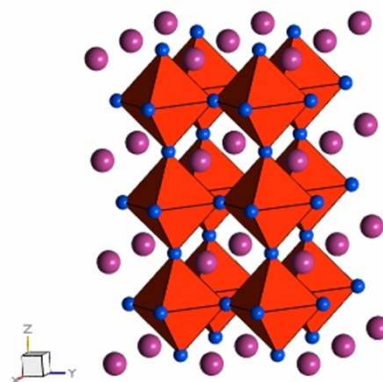
Reminder

- ↪ Equal amount of Co⁴⁺ and Co³⁺
- ↪ Co ions in octahedral coordination
- ↪ La_{1/2}Ba_{1/2}CoO₃ (~ La_{1-x}Sr_xCoO₃)



Perovskite structure Pm3m
at RT ($a_p \times a_p \times a_p$)

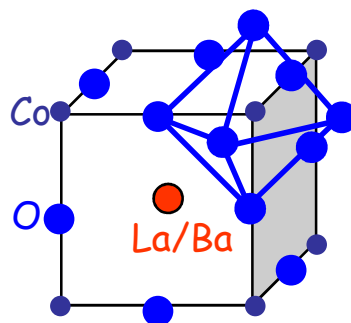
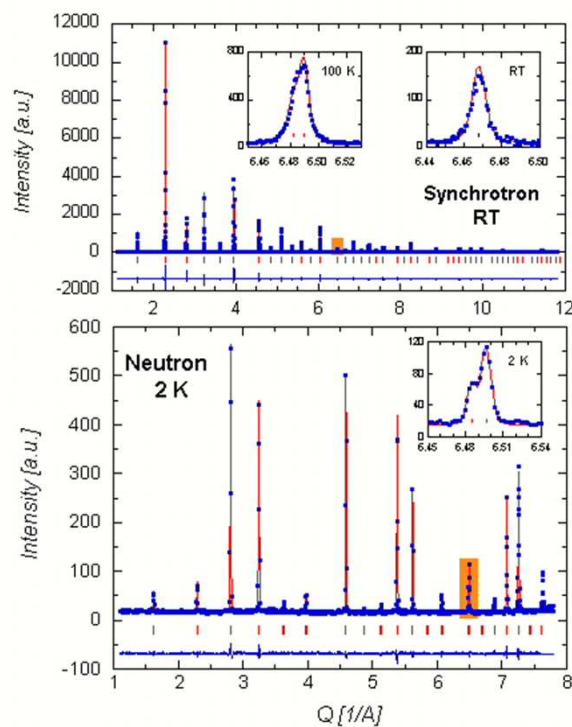
→ La and Ba statistically
distributed on the
body center site



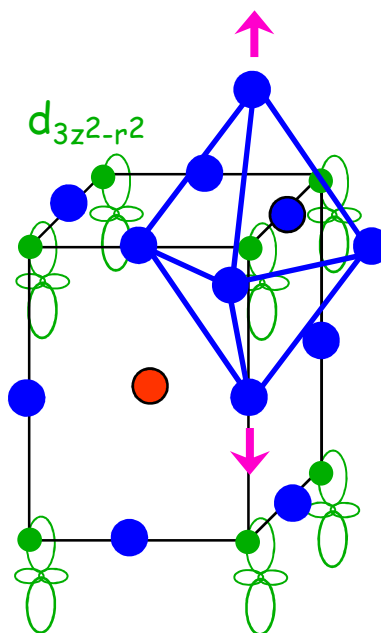
Ferromagnetism
 $T_c \sim 180$ K, $m_o(\text{Co}) = 1.9m_B$
(weak value !)

When high angular resolution is needed

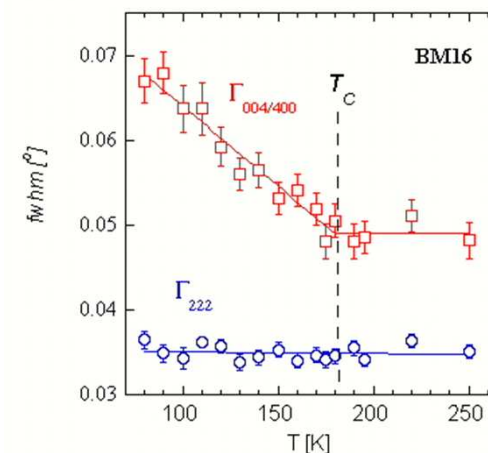
$Pm3m \rightarrow P4/mmm$
 $axaxa \quad axaxa' (a' < a)$
 $(a' - a)/a \sim 0.16\% @ 2K$



"Elongated"
octahedra



↪ Associated magnetic and structural transitions



↪ No phase segregation

Ferro-orbital ordering
of the $d_{3z^2-r^2}$

Jahn-Teller induced
ferrodistortive transition
 \rightarrow favored if IS- $Co^{3+} : |t_{2g}^5 e_g^1 \rangle$
 and/or IS- $Co^{4+} : |t_{2g}^4 e_g^1 \rangle$

Resonant X-ray scattering (often mentioned, rarely applied in powder)

EUROPHYSICS LETTERS

1 January 2001

Europhys. Lett., 53 (1), pp. 72–78 (2001)

Direct observation of 1-dimensional charge order below the first-order “metal-insulator” transition in Yb_4As_3

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Physica B 318 (2002) 284–288

PHYSICA B

www.elsevier.com/locate/physb

Charge order and crystal structure below the first-order “metal-insulator” transition in Yb_4As_3

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PHYSICAL REVIEW B 71, 075115 (2005)

Temperature dependence of the crystal structure and charge ordering in Yb_4As_3

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D. Mannix,⁶ and A. Ochial⁷¹ Swiss Light Source, Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland² European Synchrotron Radiation Facility, F-38043 Grenoble Cedex, France³ Laboratoire de Cristallographie, CNRS, F-38042 Grenoble, France⁴ Chemistry Division, Argonne National Laboratory, Argonne, Illinois-60439, USA⁵ PNC-CAT, Advanced Photon Source, Argonne National Laboratory, Argonne, Illinois 60439, USA⁶ European Synchrotron Radiation Facility, BP220, F-38043 Grenoble Cedex, France⁷ Center for Low Temperature Science, Tohoku University, Sendai, 980-8578, Japan

(Received 29 March 2004; revised manuscript received 24 November 2004; published 22 February 2005)

Yb_4As_3 : charge localization at metal-insulator transition $T_{\text{MI}}=290 \text{ K}$

I-43d \rightarrow R3c transition

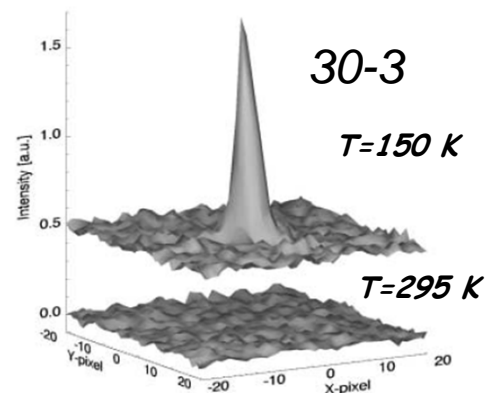
Above T_{MI} : Yb in 2.25 valence state (assuming As^{3-})

Below T_{MI} : charge ordering, two Yb species (Yb^{3+} and Yb^{4+})

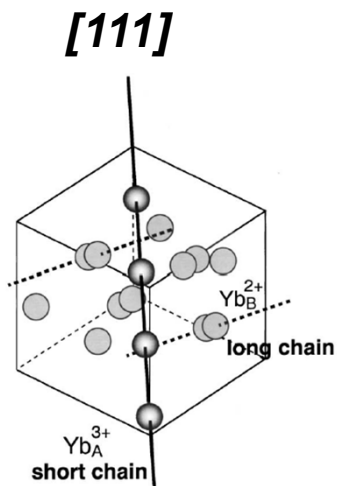
$$F(\mathbf{Q}_{\text{hkl}}) = \sum_j f_j e^{i\mathbf{Q}_{\text{hkl}} \cdot \mathbf{r}_j}, f_j = f_o + f'(E) + i \cdot f''(E)$$

$$\sum_A \exp[i\mathbf{Q} \cdot \mathbf{R}_A] = - \sum_B \exp[i\mathbf{Q} \cdot \mathbf{R}_B]$$

$F \neq 0$ if $f', f'' \neq 0$



Single crystal
CCD detector



*Intensity have to be
corrected for absorption*

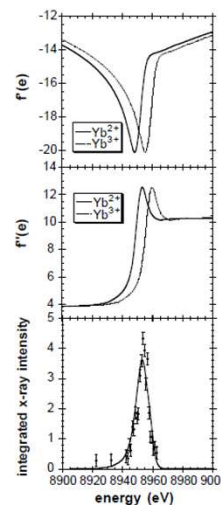


Fig. 3

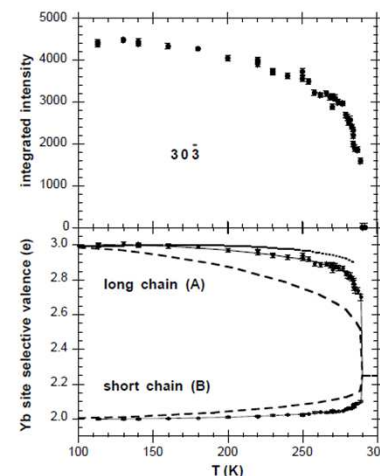
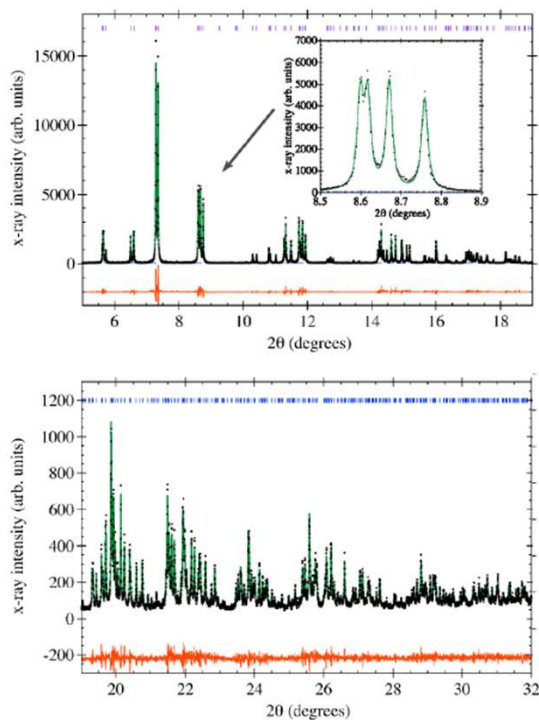


Fig. 4

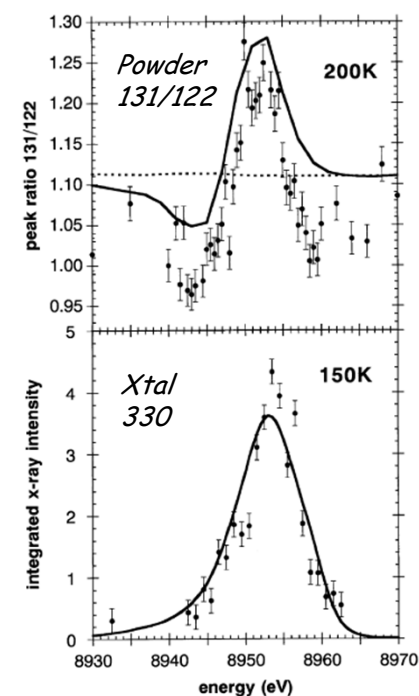
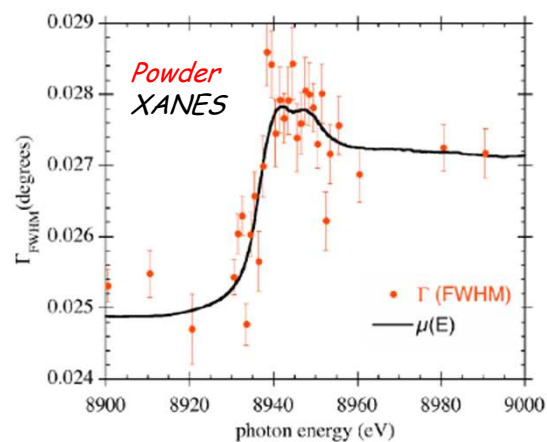
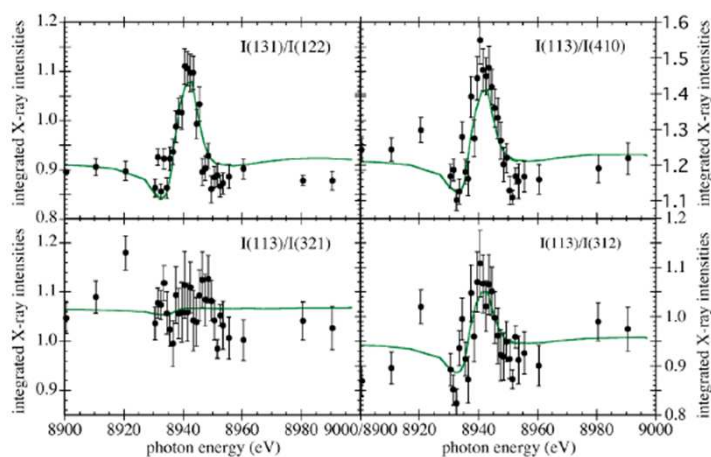
Fig. 3 – Energy dependence of the real (f') (top), and the imaginary (f'') (middle) parts of the scattering factor of Yb^{2+} and Yb^{3+} , and of the 303 reflection together with a fit to $\alpha|f_{\text{Yb}^{2+}} - f_{\text{Yb}^{3+}}|^2$, in the vicinity of the Yb L_3 -absorption edge in Yb_4As_3 .

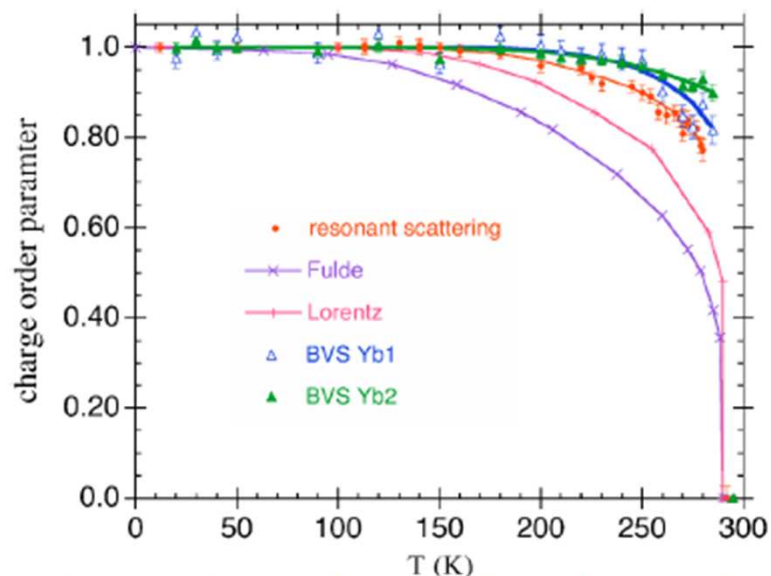
What about using powder

No need to correct for absorption since intensity ratio



$T=180\text{ K}$, $R3c$
 $a=12.503\text{ \AA}$, $c=7.496\text{ \AA}$



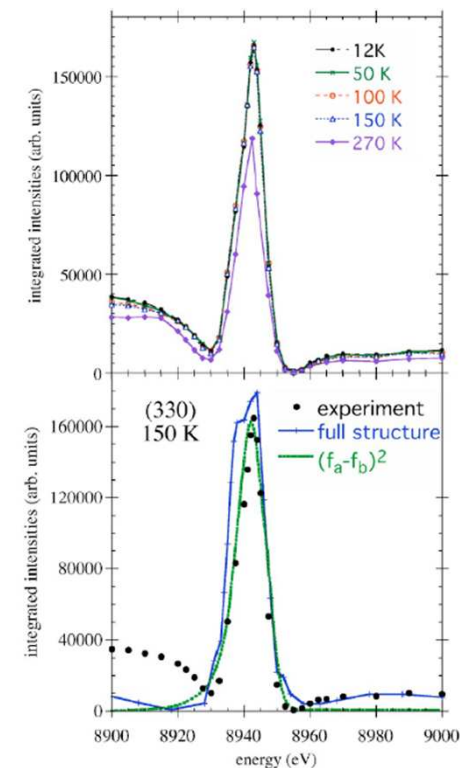


the expected valence change is weak as well. A convenient way to correlate the bond-length changes with the valence changes is to calculate the “bond valence sum” (BVS), defined as

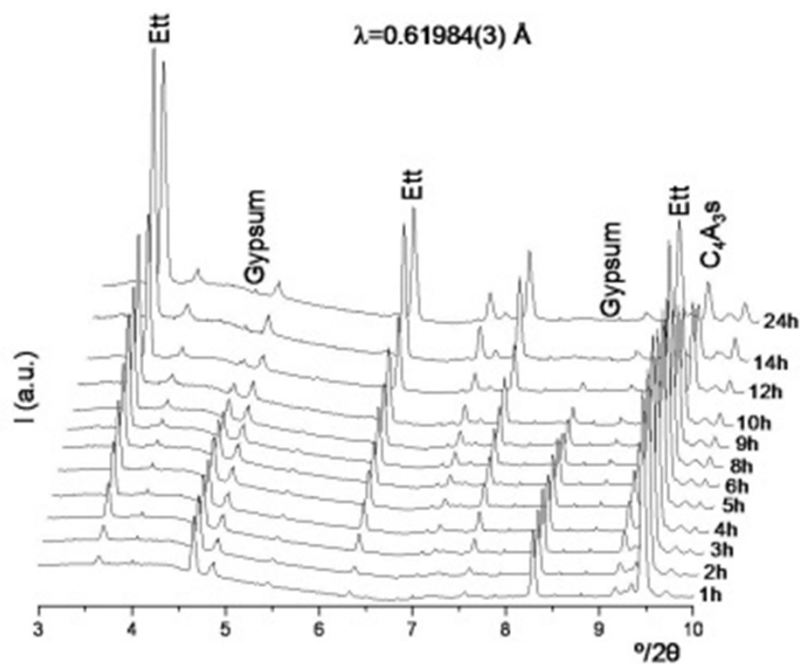
$$V_i = \sum_j \exp[(R_{ij} - d_{ij})/b],$$

*Bon valence sum
formulation*

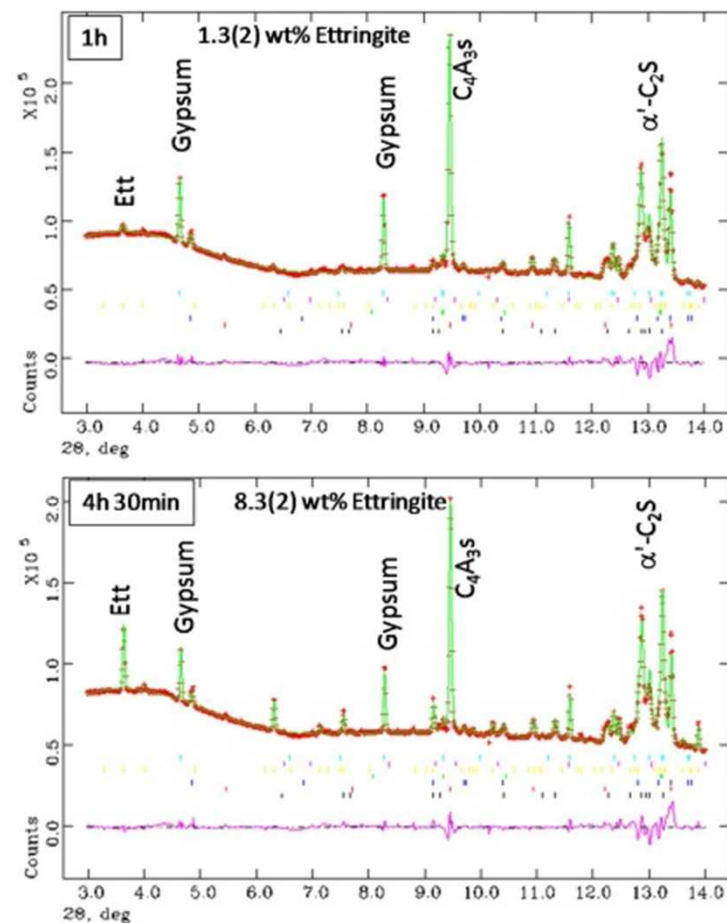
with $b=0.37$ as the general parameter, R_{ij} as the bond length between the i th and j th ion and $d_{ij}=2.59 \text{ \AA}$ ^{52,53} the parameter for the Yb—As bond, which gives the correct Yb³⁺ valences for Yb in YbAs. To correct for thermal expansion, we rescale the bond lengths with the cube root of the temperature-dependent cell volume. The obtained low temperatures value for Yb(1) and Yb(2) are 2.27 and 1.81, respectively, which deviate significantly from the expected 3+ and 2+ values. Finally, the temperature-dependent valence is obtained by scaling the corresponding BVS by the expected valence changes of $3-2.25=0.75 e$ and $2.25-2=0.25 e$, for the



First PD User experiment (October 2, 2012), Malaga University Group
 Hydration of eco cement (Miguel A.G. Aranda^{L,±}, Ángeles G. De la Torre^{*L})



Gema Álvarez-Pinazo^L, et al submitted



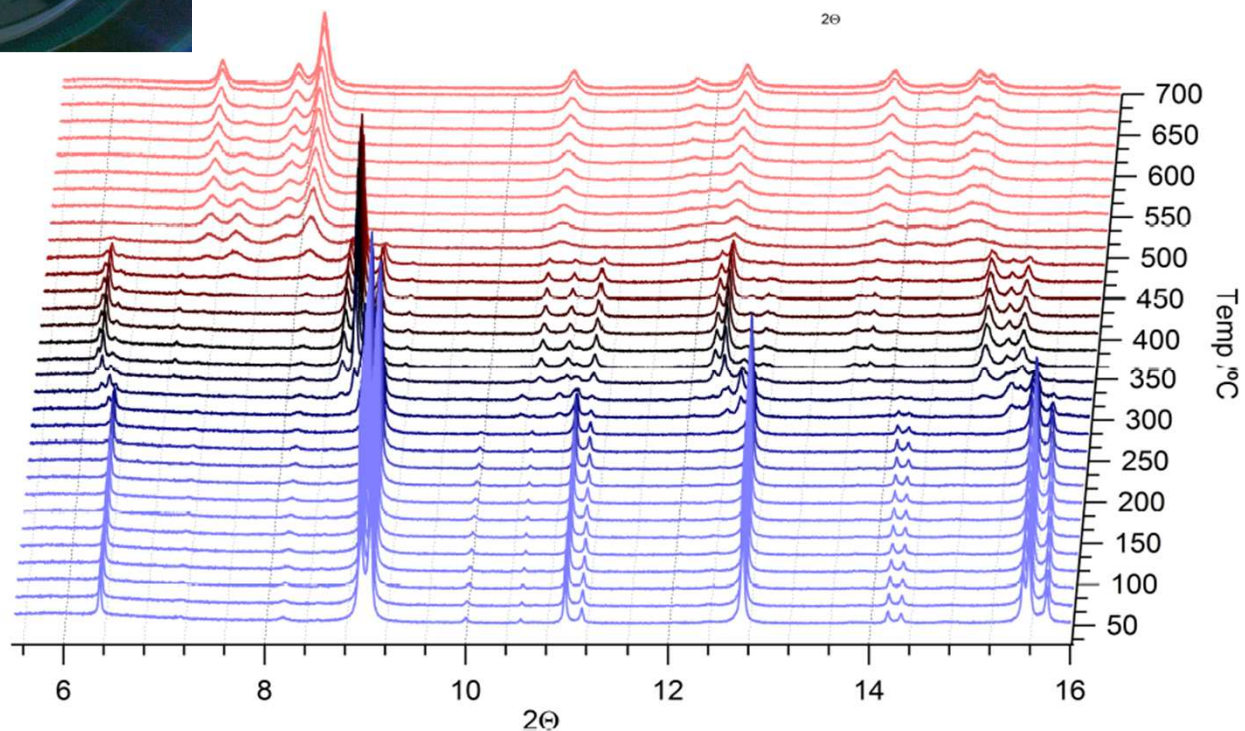
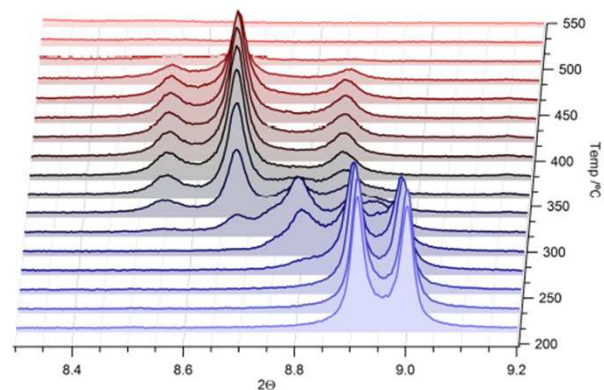
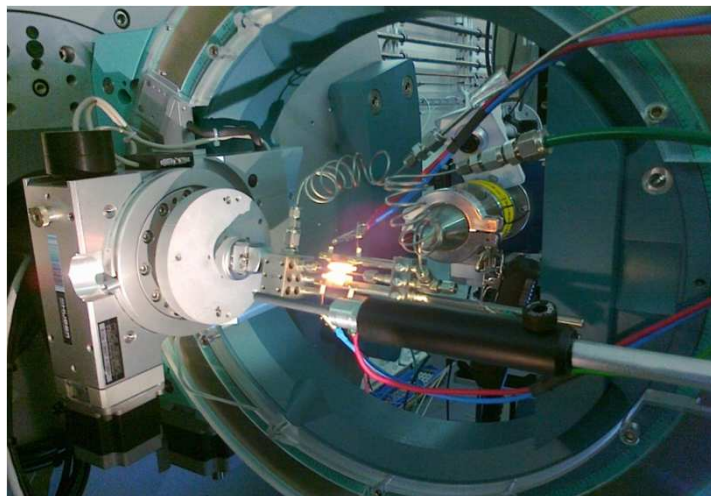
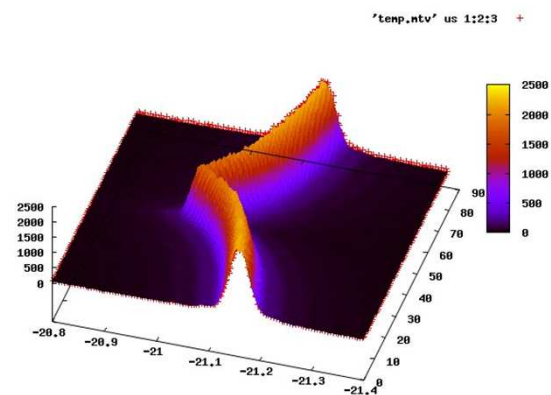
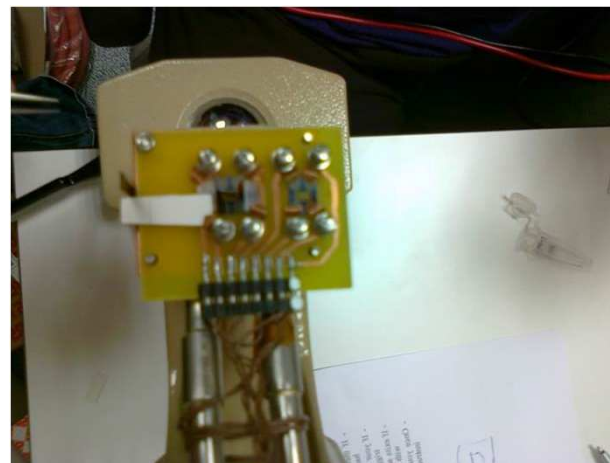


Fig 2 Reduction of LaCoO_3 to $\text{La}_2\text{O}_3 + \text{Co}$ under 5% H/He gas flow. The transition is passing through a brownmillerite intermediate. Each temperature scan was collected in 8 minutes.

When high throughput is useful



*Pt peak on thin film heated up to 800 C
And down to RT at 15deg/sec **0.5 sec/frame***



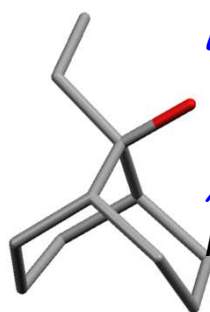
Brunelli, Fitch, Mora, Vaughan and Wright,

Structure determination by direct methods often limited by severe peaks overlap at high angle ($d \sim 1.2 \text{ \AA}$) (e.g. big unit cell, low symmetry)

Idea: Exploit anisotropic expansion of cell parameters with temperature

Severe peaks overlap at T_1 resolved at T_2

*-> 4 molecules and 48 (C+O) non-hydrogen atoms in the asymmetric unit
(attempt on BM16-PD failed, record 32 (C+O) non-hydrogen atoms in the asymmetric unit)*



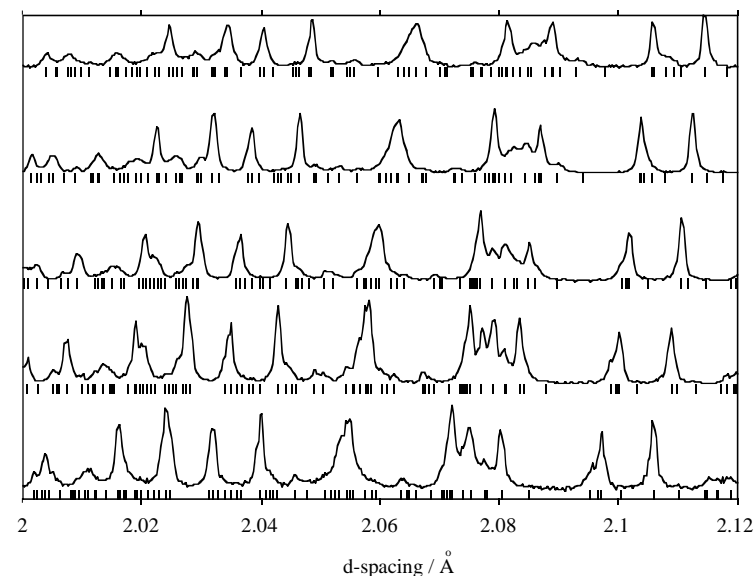
*Example: Small molecule crystallography
9-ethyl-bicyclo[3.3.1]nona-9-ol*

$C_{11}H_{20}O$

Amorph at high T , well cristallized below RT

Previously indexed in $Pbca$

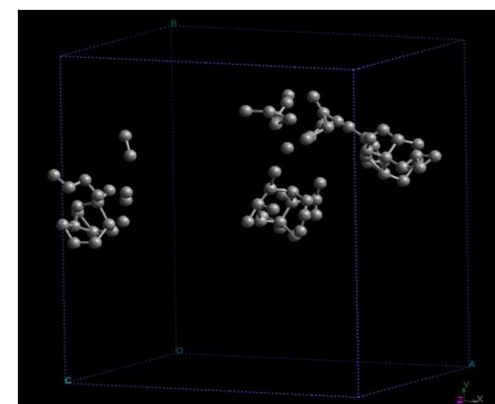
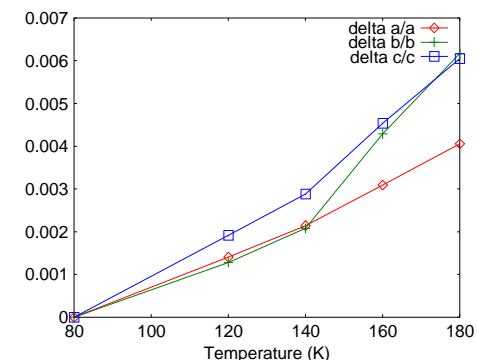
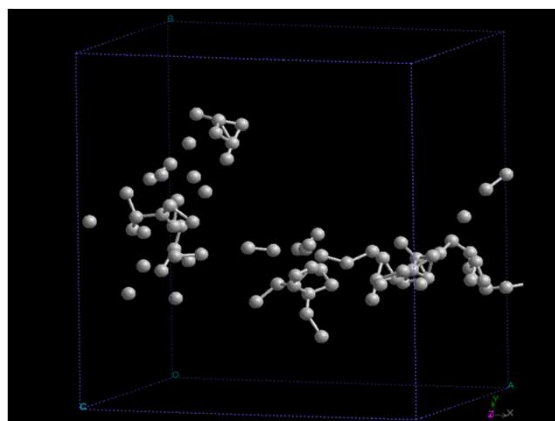
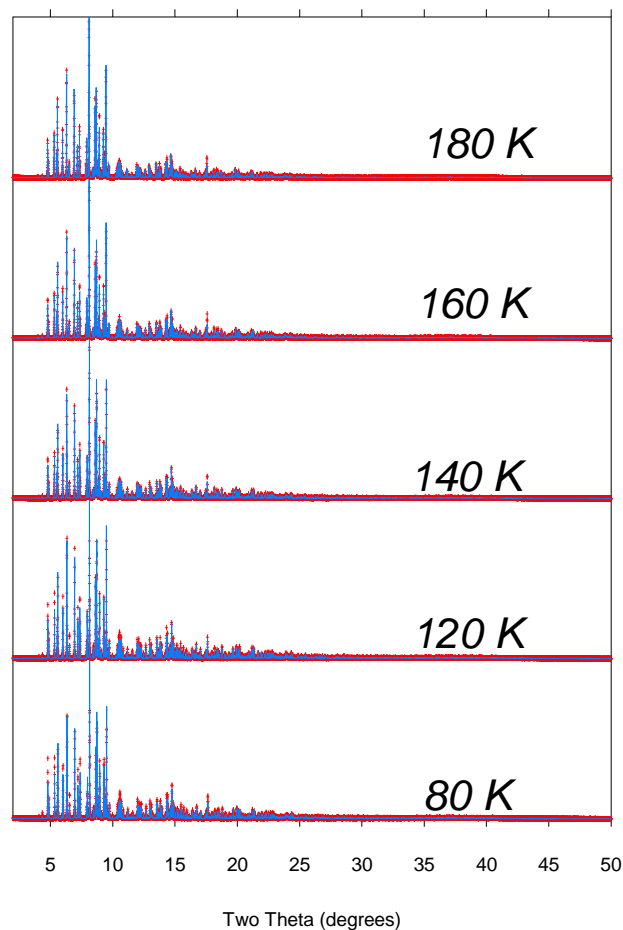
$a=21.02 \text{ \AA}$, $b=23.04 \text{ \AA}$, $c=16.38 \text{ \AA}$



When high angular resolution is needed,...but temperature helps

*3 molecules identified from initial shelxs solution (direct methods)
when 5 patterns were used.*

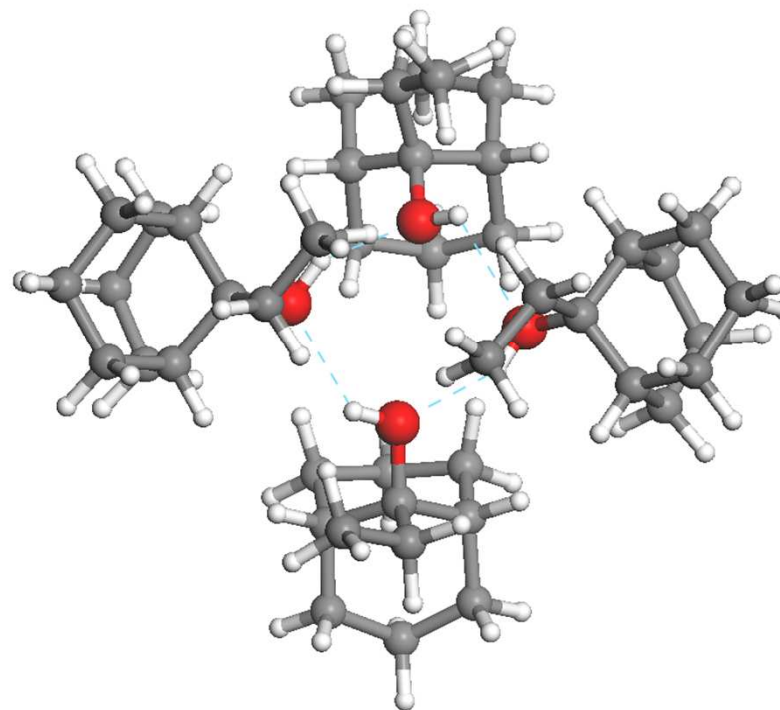
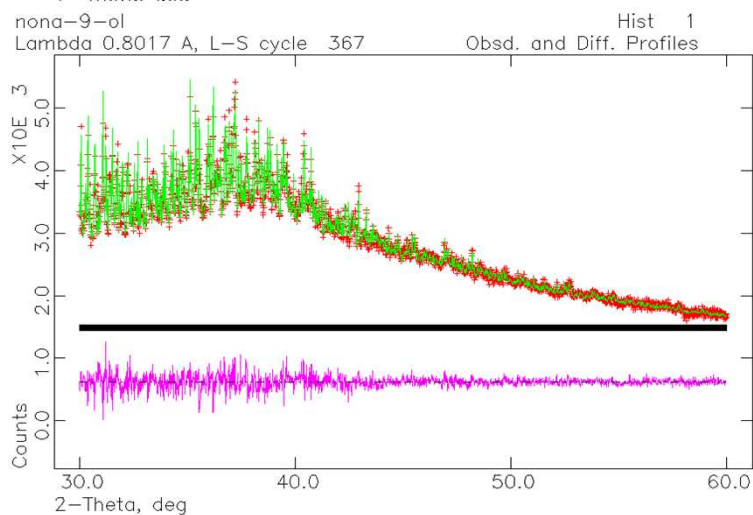
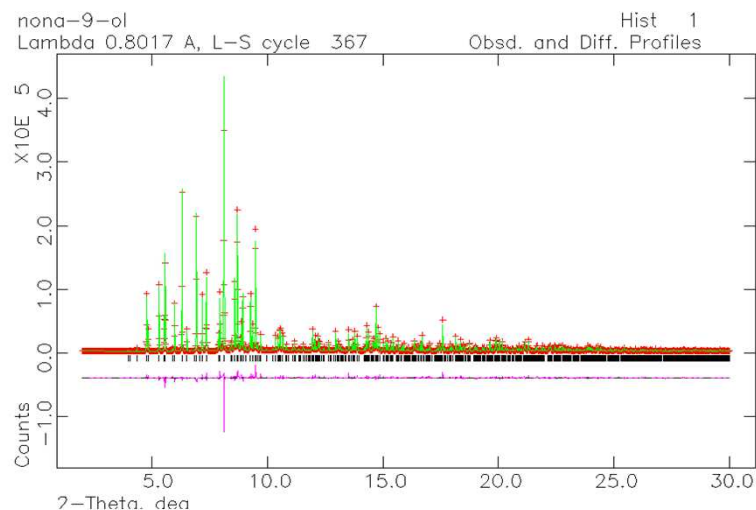
First attempt with 2 datasets (80K, 120K) and feeding intensities to shelxs direct methods failed



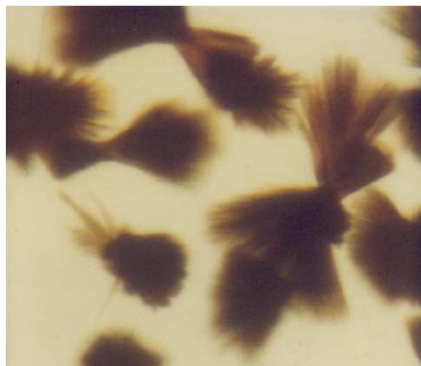
*Structure completed by Fourier recycling
against the extracted intensities from the 5
patterns*

*Refined model against 140 K data
with stereochemical restraints (GSAS)*

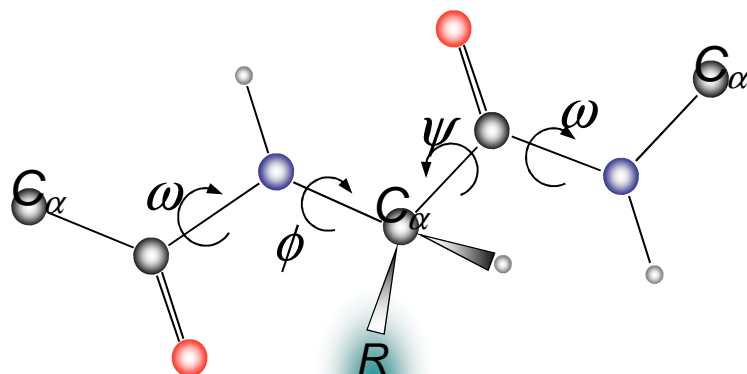
$R_{wp}=5.6\%$ $R_p=4.7\%$



Synchrotron : Powder Methods + Proteins?



Myoglobin from horse skeletal muscle
(Aldrich) recrystallised in 80% sat. $(\text{NH}_4)_2\text{SO}_4$



R. B. Von Dreele, "Combined Rietveld and stereochemical restraint refinement of a protein crystal structure." *J. Appl. Cryst.* **32**, 1084-1089 (1999).

R. B. Von Dreele, "Binding of N-acetylglucosamine to chicken egg lysozyme: a powder diffraction study." *Acta Cryst. D* **57**, 1836-1842 (2001).

J. Wright, I Margiolaki, ESRF

*Samples easier to produce... but,
Large unit cells*

Severe peak overlap

Many degrees of freedom in the structure

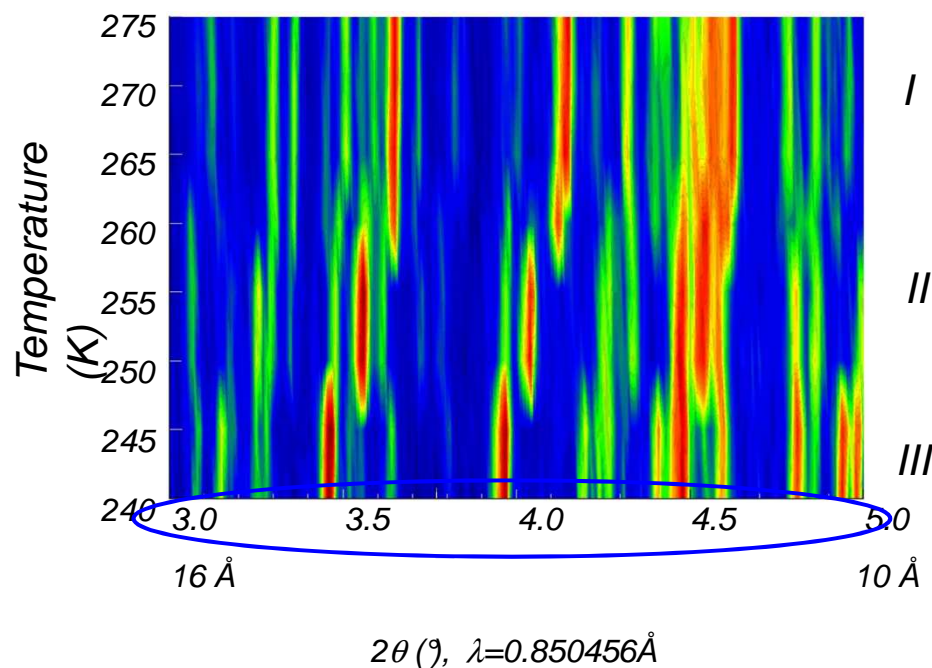
Fundamental differences to small molecule crystallography:

- Well known geometric constraints (polypeptide chains)
- only ϕ , ψ and sidechain conformations which are unknown

Improving the data quality...

*Lower temperatures for better data quality
(B factor & radiation damage)*

*Exploit anisotropic thermal expansion
to alleviate peak overlap*

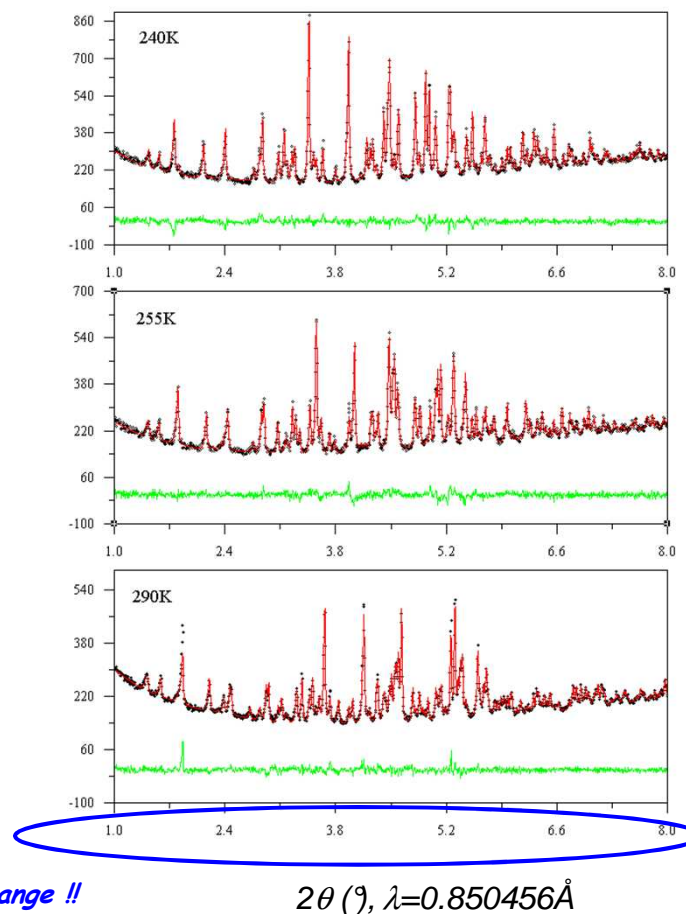


!! Small angular range !!

Combined fit for the 3 phases

*Multipattern fit with the same
intensities for each*

*Cell parameters and peak shape different for
each pattern (indexed with DICVOL)*



When small spot size is useful

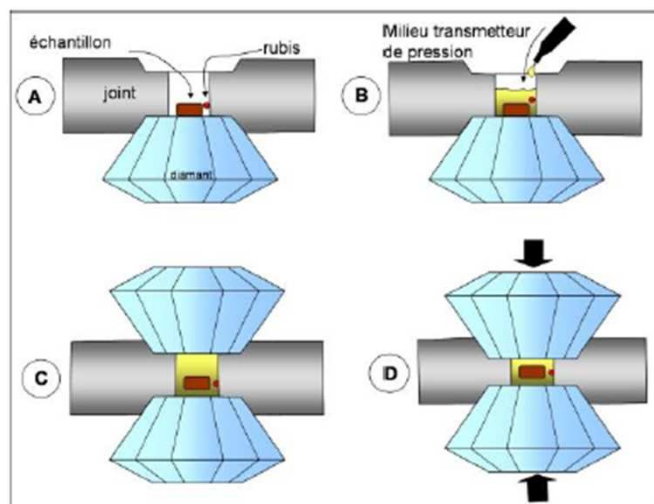
Why using small spot size ?

To illuminate a small part of the sample (spatial resolution)

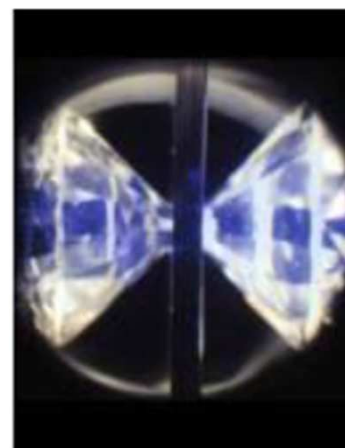
*-> Useful for mesh mapping, e.g. cultural heritage samples
(e.g. determine from crystal structure different pigments
used in paintings)*

-> strain scanning experiments (engineering)

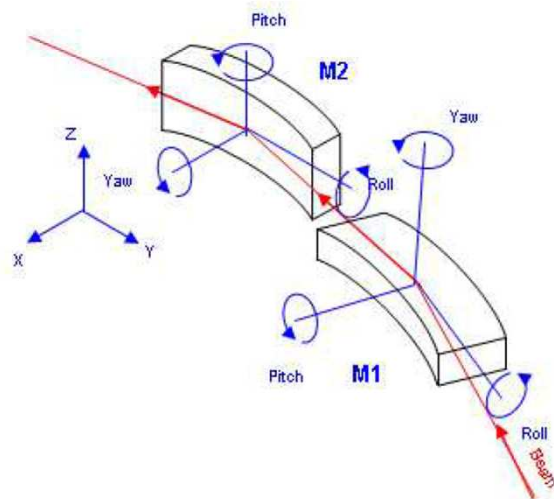
*-> unavoidable in High Pressure experiments
(small samples -> high pressure, indeed $\text{Pressure} = \text{Force} / \text{Surface}$)*



<http://www.ens-lyon.fr/LST/HP>



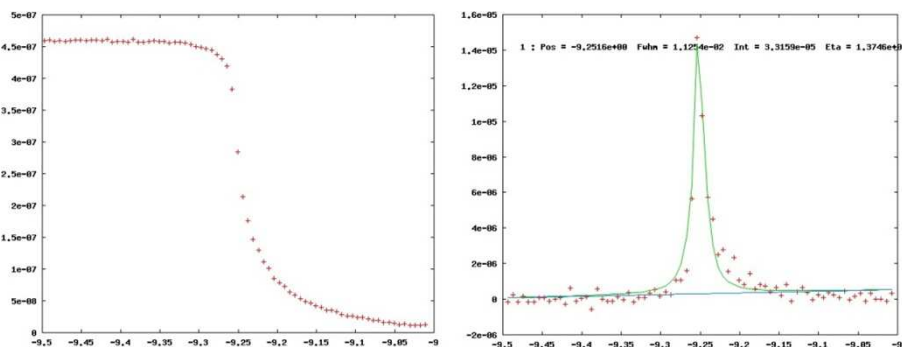
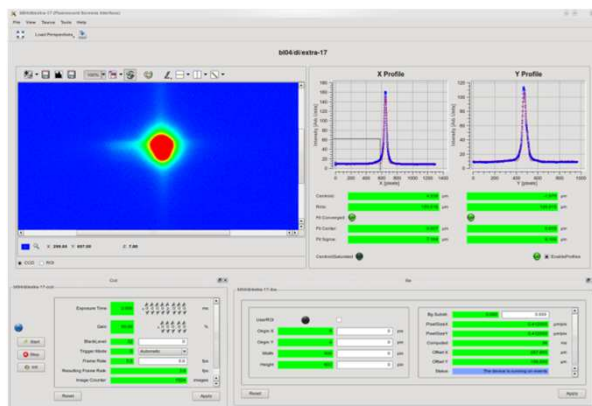
← Incoming beam,
High energy to pass
Through the diamonds



Kirkpatrick-Baez focusing optics - Irrelec (ESRF licence), France

- Elliptical bendable,
- Grating [W/Si] multilayer :
 Vert : $p/q = 29.9/1.1$, $[W/Si]_{100}$, $d \sim 3\text{nm}$, slope error $< 1\ \mu\text{rad}$
 Hor : $p/q = 30.3/0.7$, $[W/Si]_{150}$, $d \sim 2.6\text{nm}$, slope error $< 1.5\ \mu\text{rad}$
- Focus spot of $12\ \mu\text{m}$ (H) \times $7\ \mu\text{m}$ (V) at sample position over 20-50 keV

KB optimized to focus at the HP station, but designed to focus at the PD station as well



-> 10 X 10 μm beam

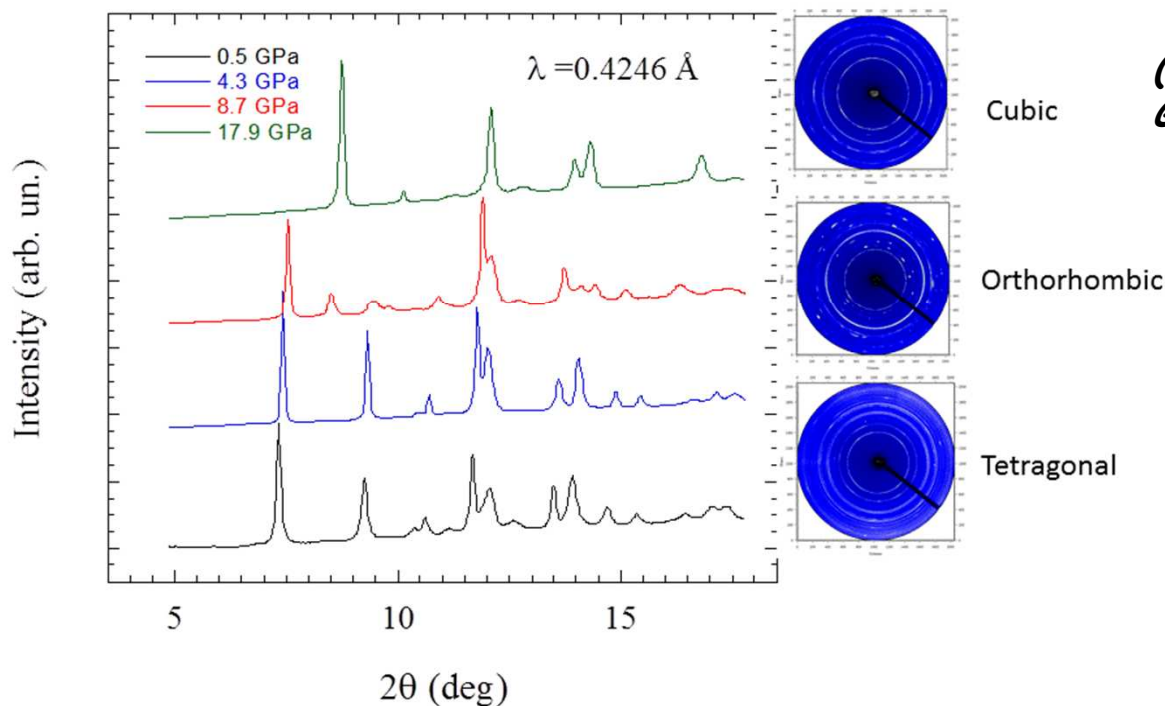
First HP User experiment (June 27, 2012), Cantabria University Group

Pressure-induced structural-phase transition in $[(C_nH_{2n+1})_4N]FeX_4$ ($n=1,4$; $X: Cl, Br$) associated with changes of Fe coordination.

Powder diffraction patterns versus pressure

Pressure-induced phase transition sequence:

Tetragonal \rightarrow Orthorhombic \rightarrow Cubic



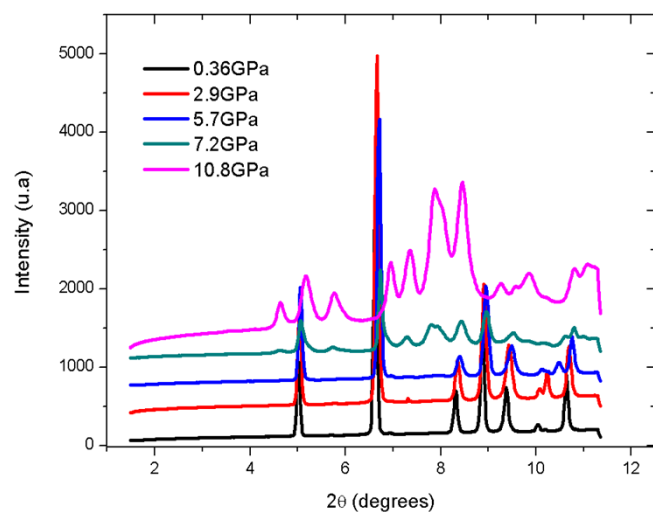
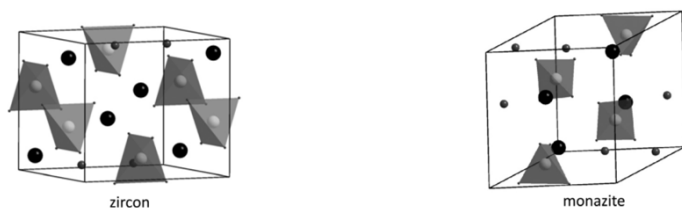
2D raw data to 1D patterns

using A Hammersley's FIT2D

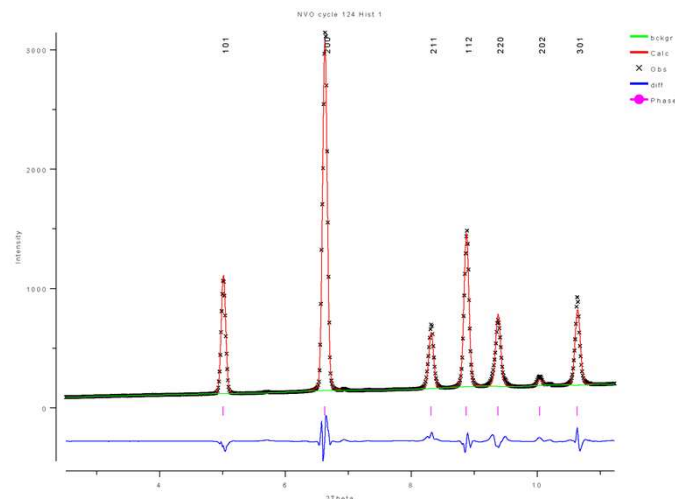
<http://www.esrf.eu/computing/scientific/FIT2D>

(evaluate possibility to use GSAS2 package)

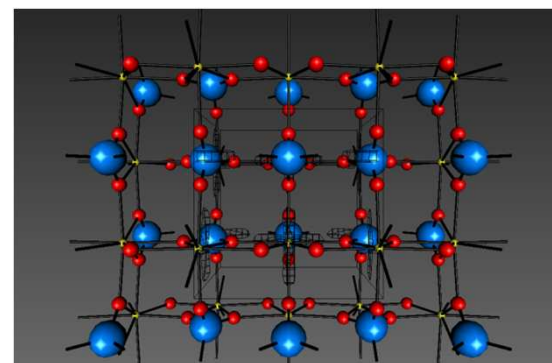
Zircon to monazite phase transition in NdVO₄ (I41/amd, Z=4) > Monazite (P21/n, Z=4)



Evolution of the XRD patterns of NdVO₄ as a function of pressure

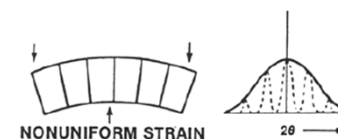
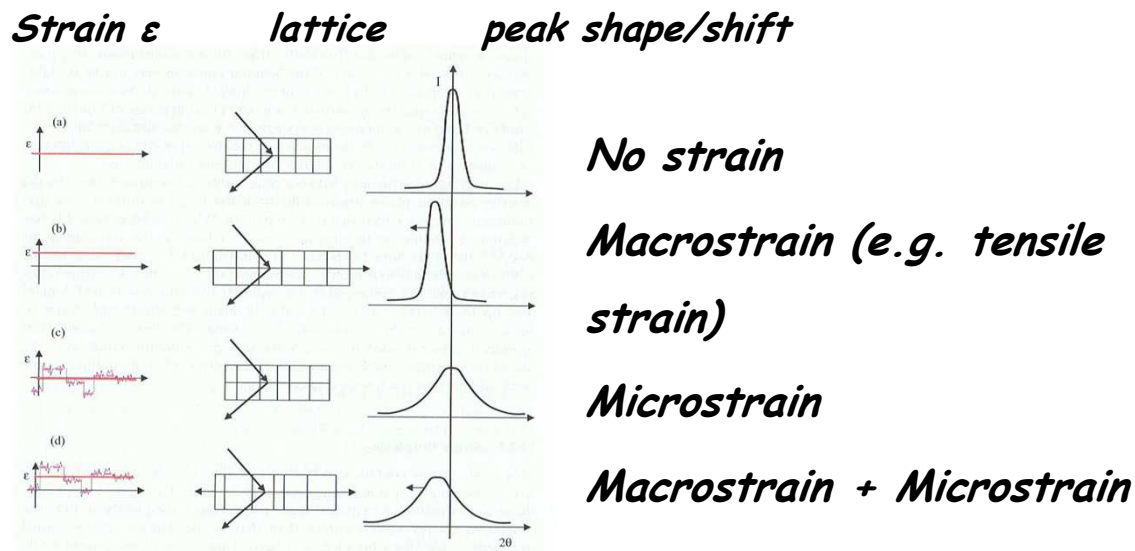


**3D difference
Fourier plot (GSAS)**



Why high resolution and high accuracy ?

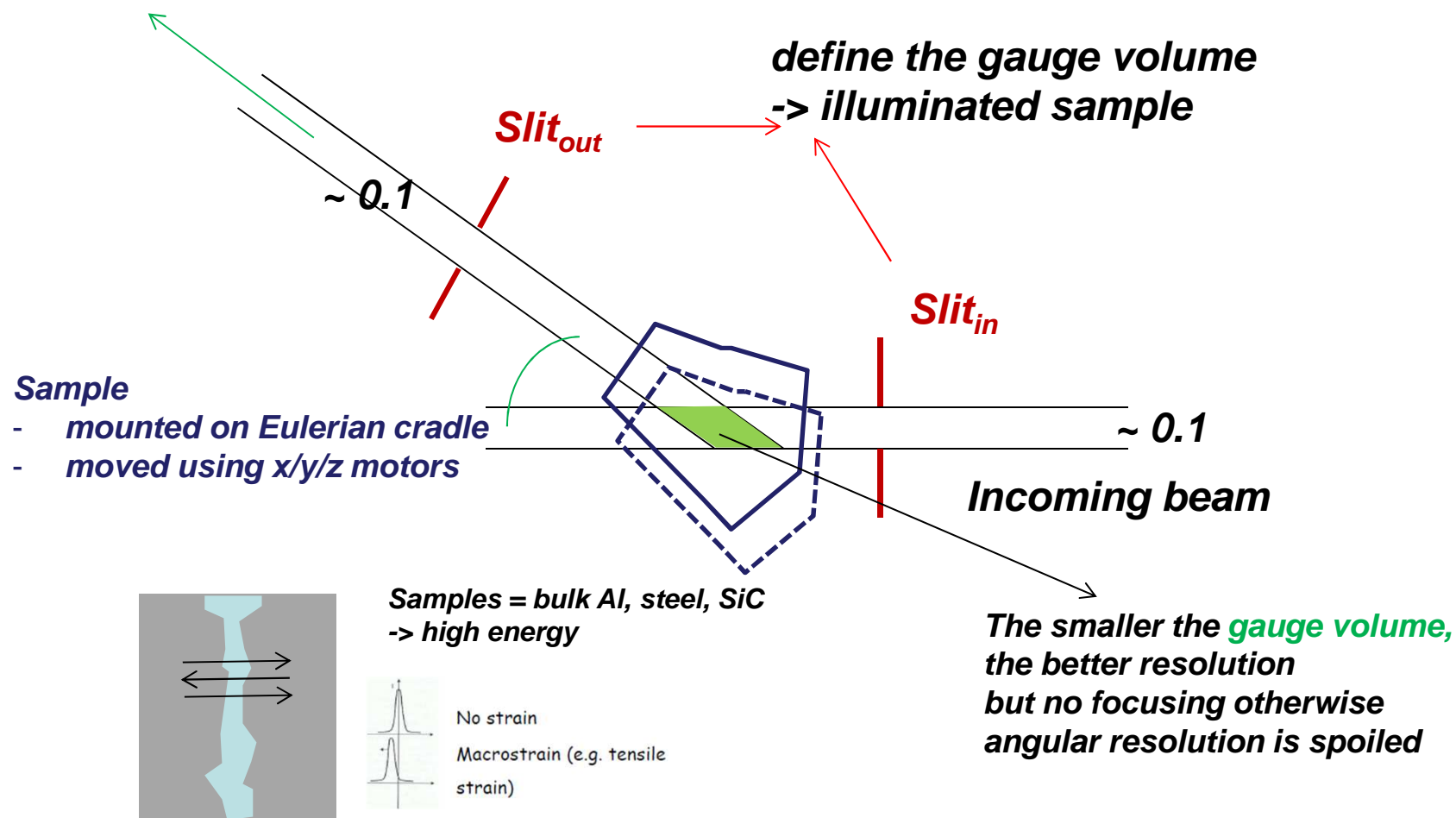
Sometimes the change in position are very tiny (e.g. change in temperature)



In a 'normal' powder without external forces no macrostrain is seen.

However, microstrain is important e.g in Metallurgy

To MAD central channel

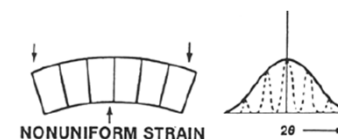
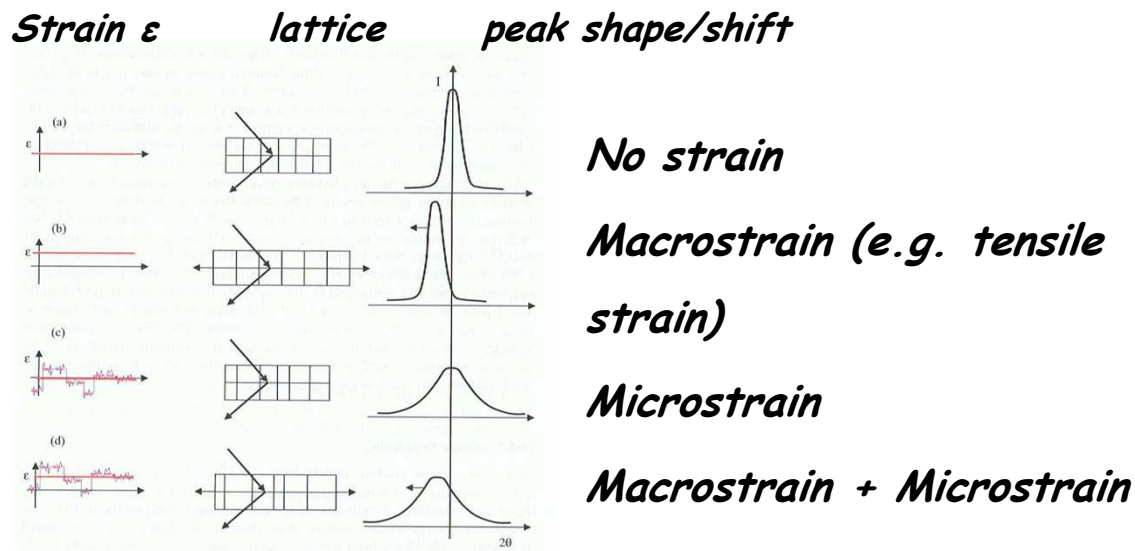


2 pieces welded by various techniques

-> goal is to have compressive strains at welding points

Why high resolution and high accuracy ?

Sometimes the change in position are very tiny (e.g. change in temperature)



In a 'normal' powder without external forces no macrostrain is seen.

However, microstrain is important e.g in Metallurgy

$\text{Pr}_{0.5}\text{Ca}_{0.5}\text{Mn}_{0.97}\text{Ga}_{0.03}\text{O}_3$, a strongly strained system due to the coexistence of two orbital ordered phases at low temperature

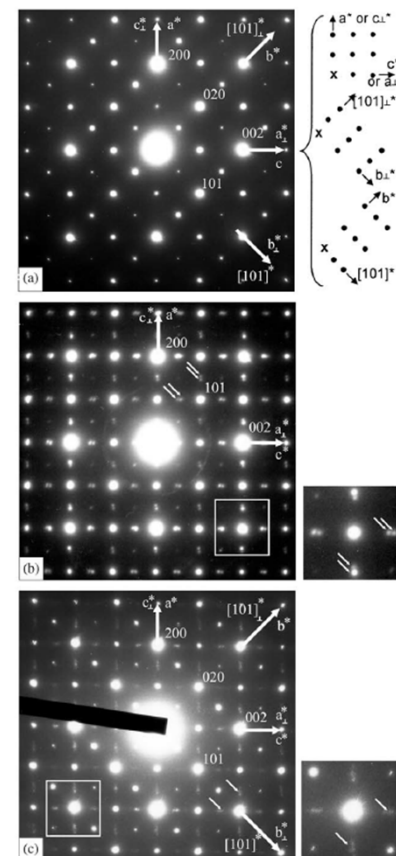
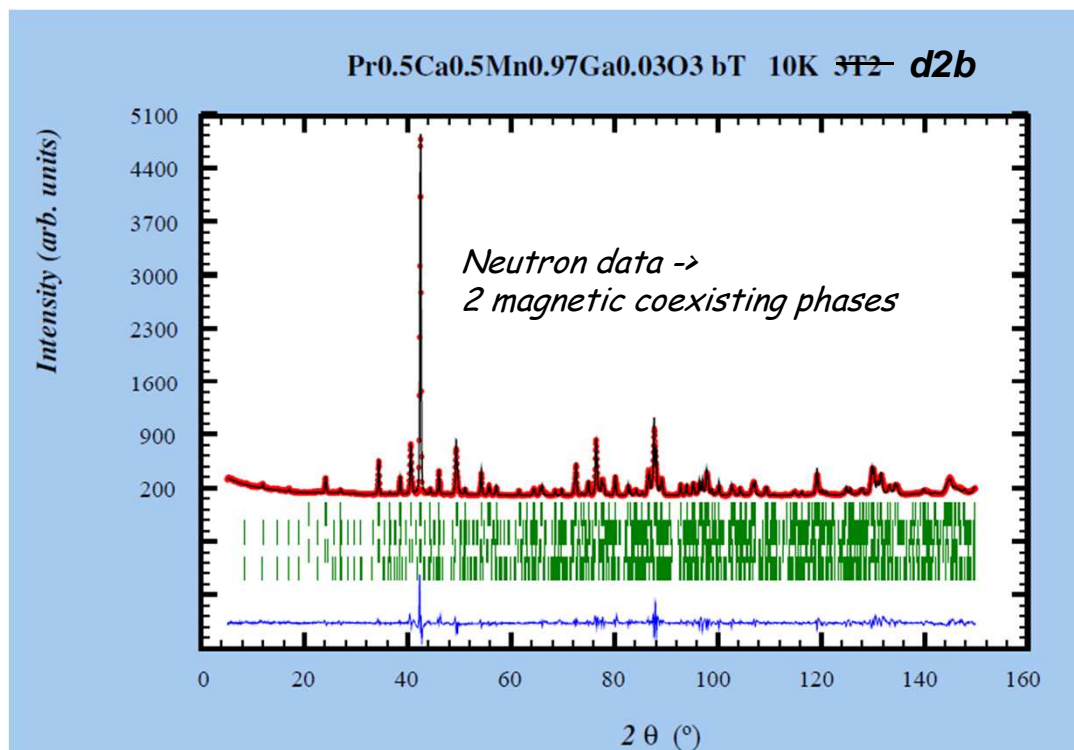
C. Yaicle^{a,*}, F. Fauth^b, C. Martin^a, R. Retoux^a, Z. Jirak^c, M. Hervieu^a,
B. Raveau^a, A. Maignan^a

^aLaboratoire CRISMAT, UMR 6508 CNRS ENSICAEN, 6 bd Maréchal Juin, 14050 CAEN Cedex, France

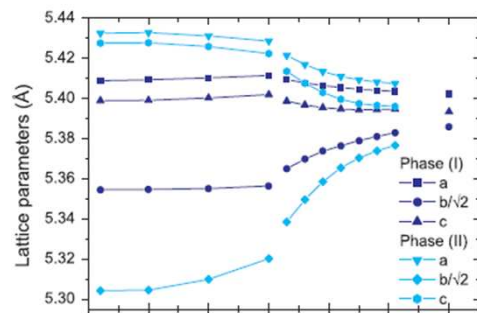
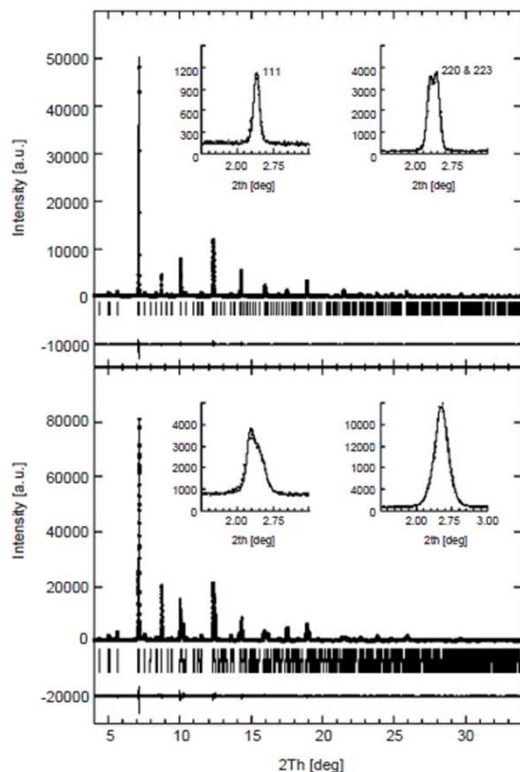
^bESRF, BP 220, 6 Rue Jules Horowitz, F-38043 Grenoble, France

^cInstitute of Physics ACSR, Cukrovarnicka 10, 162 53 Prague 6, Czech Republic

Electron diffraction data →
2 crystal coexisting phases



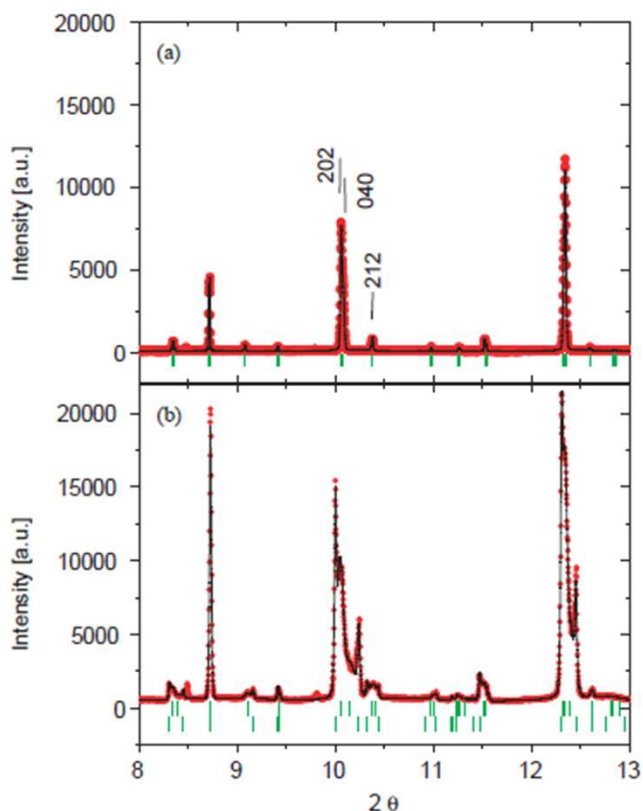
When high angular resolution is needed...for microstructural effects



Anisotropic line broadening, using Stephens formalism

$$\Gamma^S(hkl) = \sqrt{\sigma^2(1/d_{hkl}^2) \tan(\theta)/(1/d_{hkl}^2)},$$

$$\sigma^2(1/d_{hkl}^2) = \sum_{\substack{i,j \\ \alpha_i=a,b,c,\gamma,\beta,\gamma}} \langle (\alpha_i - \langle \alpha_i \rangle)(\alpha_j - \langle \alpha_j \rangle) \rangle \\ \times \frac{\partial(1/d_{hkl}^2)}{\partial \alpha_i} \frac{\partial(1/d_{hkl}^2)}{\partial \alpha_j} \\ = \sum_{\substack{HKL \\ H+K+L=4}} S_{HKL} h^H k^K l^L.$$



The anisotropic line broadening reflects the local structural distortions observed by electron diffraction/ microscopy

They are as many types of synchrotron powder diffraction experiments as users' ideas

Make use of wide range of energies, high resolution, high throughput, small beam size, heavy duty diffractometers, big space around sample position, low/high temperatures, high pressures, gas flow,...

***Make use of available complementary techniques :
Neutron powder diffraction, electron diffraction/microscopy,
transport properties (resistivity, DSC), chemical analysis***

***Sorry for not having mentioned Pair Distribution Function
type data acquisition***